1	Magnetic Local Time-resolved Examination of Radiation Belt Dynamics
2	during High Speed Solar Wind Speed-Triggered Substorm Clusters
3	Craig J. Rodger
4	Department of Physics, University of Otago, Dunedin, New Zealand
5	Drew L. Turner
6	Space Sciences Department, Aerospace Corporation, El Segundo, CA, United States
7	Mark A. Clilverd
8	British Antarctic Survey (NERC), Cambridge, United Kingdom
9	Aaron T. Hendry
10	Department of Space Physics, Institute of Atmospheric Physics, Prague, Czechia
10 11	Department of Space Physics, Institute of Atmospheric Physics, Prague, Czechia
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10 11 12 13	Department of Space Physics, Institute of Atmospheric Physics, Prague, Czechia Main point # 1: Superposed epoch analysis around clusters of substorms show consistent radiation belt dynamical responses to mild geomagnetic disturbances.
10 11 12 13 14	Department of Space Physics, Institute of Atmospheric Physics, Prague, Czechia Main point # 1: Superposed epoch analysis around clusters of substorms show consistent radiation belt dynamical responses to mild geomagnetic disturbances. Main point # 2: Magnetopause shadowing produces proton and electron loss over a wide
10 11 12 13 14 15	Department of Space Physics, Institute of Atmospheric Physics, Prague, Czechia Main point # 1: Superposed epoch analysis around clusters of substorms show consistent radiation belt dynamical responses to mild geomagnetic disturbances. Main point # 2: Magnetopause shadowing produces proton and electron loss over a wide range of <i>L</i> -shells, eventually modified by substorm flux enhancements.
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23 different MLT, such as magnetopause shadowing, plasma wave activity, and substorm

injections. Analysis shows that magnetopause shadowing produces clear loss in proton and 24 electron populations over a wide range of L-shells, initially on the dayside, which interact 25 with nightside substorm-generated flux enhancements following charge-dependent drift 26 directions. Inner magnetospheric injections recently identified as an important source of 27 10's to 100's keV electrons at low L (L < 3), occurring during similar solar wind-driving 28 conditions as recurrent substorms, show similar but more enhanced geomagnetic AU-index 29 signatures. Two-fold increases in substorm occurrence at the time of the sudden particle 30 enhancements at low L shells (SPELLS), suggests a common linkage. 31

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Plain English Summary: The magnetic field of the Earth is filled with high-energy 33 particles, primarily electrons and protons, forming the Van Allen radiation belts. Over the 34 years it has become obvious that the number of trapped high-energy electrons changes 35 rapidly, and in complex ways. We know that multiple different processes are involved to 36 produce such dynamic changes, which include energization, transport, and loss. Over the 37 last ~7 years flagship science missions have been launched by multiple space agencies to 38 better understand the complex dynamics. However, these involve only 1 or 2 highly 39 instrumented spacecraft - these make extremely high quality measurements, but are limited 40 by their inability to be in multiple places at the same time. Nonetheless, signatures of 41 several different processes have been identified and described. In the current study we use a 42 constellation of spacecraft with more limited instrumentation than the flagship missions. 43 This has allowed us to clarify the typical dynamical processes affecting radiation belt 44 particles. In particular, the ability to measure simultaneously at multiple locations, plus 45 statistical averaging, allows us to show clear evidence of the loss process termed 46 magnetopause shadowing. We also cast light on a previously mysterious process of sudden 47 particle enhancements occurring deep in the belts. 48

50 1. Introduction

The temporal evolution of radiation belt electron fluxes is highly dynamic, particularly for 51 the outer radiation belt. Multiple different processes have been identified which can drive 52 electron energization (often termed acceleration), loss, or transport [see for example, Balasis 53 et al., 2016; Baker et al., 2018]. Typically the occurrence and magnitude of these processes 54 are dependent upon distance from the Earth (expressed, for example, through the L-shell), 55 particularly due to the changing cold plasma density and the strong gradients around the 56 plasmapause. At the same time the processes also depend very strongly on magnetic local 57 time (MLT). The trapped radiation belt electron flux at a given point in space at a given time 58 depends on a combination of multiple processes - in order to understand the evolution of 59 trapped flux it is necessary to understand the MLT-dependent dynamical processes in some 60 detail. Typically, MLT-dependence is averaged over processes that take significantly longer 61 than the electron drift period, and as such have not been clearly seen in many experimental 62 studies. 63

There are multiple examples of L and MLT-dependent activity influencing radiation belt 64 electron flux evolution on timescales faster than the electron drift period. One is 65 magnetopause shadowing on the dayside [e.g., Turner et al, 2012; Yu et al., 2013] where 66 electrons drifting around the Earth encounter the magnetopause and are lost into the solar 67 wind. Another is the strong MLT variation seen in plasma wave activity [e.g., Figure 7 of 68 Summers et al., 1998], which is likely responsible for the L and MLT variation seen in 69 precipitation into the atmosphere [e.g., Carson et al., 2012; Douma et al., 2017; van de Kamp 70 et al., 2018]. A third example is magnetospheric substorms, short-lived reconfigurations of 71 the geomagnetic field where energetic particles are injected into the inner magnetosphere 72 close to magnetic midnight [Akasofu, 1981; Cresswell-Moorcock et al., 2013]. 73

In many cases we have a good physical understanding of how the L and MLT-dependent 74 activity will drive changes in radiation belt fluxes; however, it is not always easy to observe 75 these dynamical changes occurring in-situ and discriminate between the actions of individual 76 processes. One strong reason for this is the rapid drift time of relativistic electrons. The drift 77 period of a trapped 1 MeV electron at L=5 is ~13.5 min (calculated through expressions in 78 *Walt* [1994]). Even a dramatic change occurring in one MLT region will rapidly drift through 79 all other sectors making it hard to determine where it originated from; such rapid drift rates 80 compared to the time resolution of the analysis also mean MLT-dependent impacts are 81 rapidly "smeared" around the Earth, especially considering spacecraft revisit periods of ~1 82 hour to a day, or more. In contrast, the drift period of a trapped 150 keV electron at L=5 is 83 almost 70 min. In the current study we use electrons with energies of hundreds of keV, which 84 drift much more slowly than MeV electrons and hence individual processes can be 85 distinguished using a network of Low Earth Orbit (LEO) satellites with ~100 minute orbital 86 periods. 87

In the last decade or so, our understanding of the radiation belts has markedly increased, in 88 large part due to flagship space missions (examples being the Van Allen Probes and Arase), 89 along with the concentrated scientific attention such large scale activity attracts. However, the 90 cost of such high-quality scientific platforms limits the number which will operate at any 91 given time. In recent years there has been a wealth of exceptional in-situ observations 92 deepening our understanding [e.g., Jaynes et al., 2015; Aseev et al., 2017; Zhao et al., 2017; 93 Kasahara et al., 2018; Xiang et al., 2018; Turner et al., 2019] but which are limited in their 94 ability to provide simultaneous MLT and L coverage. 95

This is an area in which spacecraft with more limited instrumentation in LEO can assist, as they already exist as constellations of multiple satellites simultaneously monitoring different MLTs while rapidly moving through *L*-shells. In this study we make use of observations from a constellation of polar-orbiting LEO satellites which have employed the same experimental equipment to make measurements of medium energy electrons in the radiation belts for >15 years. Using these observations as a "big data" set, we undertake superposed epoch analysis (SEA) around clusters of substorms, inner magnetospheric activity expected to produce a strong radiation belt dynamical response [*Miyoshi et al.*, 2013; *Jaynes et al.*, 2015; *Rodger et al.*, 2016]. The analysis demonstrates the high level of dynamical variation in radiation belt structure with MLT. By focusing on medium energy electrons we can discriminate between processes occurring at different MLT.

107 **2. Experimental Datasets**

108 2.1 POES SEM-2 particle observations

The Polar Orbiting Environmental Satellites (POES) are a set of LEO spacecraft (~800-850 109 km) in ~100-minute period Sun-synchronous polar orbits. Since NOAA-15 in 1998, this 110 series of spacecraft have monitored medium energy electron and proton fluxes with the 111 Medium Energy Proton and Electron Detector [Evans and Greer, 2004; Rodger et al., 2010a, 112 Rodger et al., 2010b] as part of the Space Environment Monitor (SEM-2) package [Evans 113 and Greer, 2004]. Here we focus on the trapped electron and proton flux observations from 114 the 90-degree telescopes, which are named 90eX and 90PX, where X is the channel number 115 (see [Evans and Greer, 2004; Rodger et al., 2010a] for more details). We restrict ourselves to 116 the 90-degree telescope observations as these primarily measure trapped particles [Rodger et 117 al., 2010a, Rodger et al., 2010b], the dynamics of which are the focus of the current study. 118

During the period analyzed in our study (2005-2013) an increasing number of POES spacecraft were launched, including the US NOAA-15 through to NOAA-19, and the European MetOp-1 and 2. Due to the large number of POES spacecraft, and their LEO orbits, there is very good coverage across L and MLT [e.g., Hendry et al., Fig. 1, 2016]. We have combined the observations from multiple satellites into an L and time grid of median flux values with a 0.25 L-resolution and a 15 min time resolution. This is undertaken for a series of MLT range: 0-3, 3-6, 6-9, through to 21-24 MLT. A more detailed description of the
dataset and the processing undertaken can be found in *Rodger et al.* [2010a] and *Cresswell- Moorcock et al.* [2013].

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129 2.2 Recurrent Substorm epochs

As noted above, we will undertake SEA on "clusters" of substorms, and make use of the 130 substorm lists produced by SuperMAG [Newell and Gjerloev, 2011a; Gjerloev, 2012]. We 131 follow the definition of *Newell and Gjerloev* [2011b] and their naming convention of 132 "recurrent" substorm groupings, the definition of which is described below. For the current 133 study the SuperMAG substorm list was generated on 25 Aug 2014, 18:57:19 UT. This is the 134 same "recurrent" substorm list used in *Rodger et al.* [2016], spanning 2005-2013. We make 135 use of this listing as its properties, links to solar wind and geomagnetic activity, and non-136 MLT dependent radiation belt SEA were described in detail in Rodger et al. [2016]. This 137 allows us to focus in this research letter specifically on the MLT dependent behavior of the 138 same set of events. 139

Recurrent Substorm Epoch: The epoch event time is taken as the time of the first substorm in a cluster of 2 or more substorms which are closely spaced in time, using the clustering definition of *Newell and Gjerloev* [2011b]. The start of the cluster must have >82 minutes between it and any previous events. Each subsequent substorm in the chain must be spaced \leq 82 minutes after its immediate previous neighbor. There is no restriction on the length of the recurrent substorm chain. Following *Rodger et al.* [2016], there are 2,052 unique recurrent substorm epochs.

147 **3. MLT-resolved Dynamical processes**

148 **3.1 Overview**

Rodger et al. [2016 showed that the recurrent substorm epochs begin during times of high 149 solar wind speeds when the IMF B_z turns southwards. Following the recurrent substorm 150 epochs there are enhancements in lower band whistler mode chorus and energetic electron 151 152 fluxes in the radiation belts. It is important to note that both the chorus and flux enhancements start before the zero epoch, consistent with acceleration driven by enhanced 153 magnetospheric convection driven by large-amplitude Alfvén waves in the solar wind [Lyons 154 et al., 2005]. However, there is a much stronger whistler mode chorus enhancement after the 155 zero epoch [Rodger et al., 2016], consistent with the importance of substorms providing a 156 population of chorus-producing source electrons [e.g., Jaynes et al., 2015; Simms et al., 157 2018]. 158

Unremarked in the text of the *Rodger et al.* [2016] study is the clear increase (shown in the right hand side of Figure 4 of that study) for inner belt and slot region >100 keV electrons occurring after the recurrent substorm zero epoch We argue this is consistent with what has recently been named sudden particle enhancements at low *L*-shells (SPELLS) [*Turner et al.*, 2017]. The SPELLS reports stimulated the research efforts detailed in this current work.

It is important to note that the epochs used here are representative of dynamical changes during weak geomagnetic disturbances, i.e., outside of geomagnetic storms. The median Dst at zero epoch only reaches -18 nT (c.f., typical value of -9 nT for all times). Only ~4.5% of the epochs show evidence of storm conditions (\leq -50 nT). As such our analysis is unlikely to be strongly influenced by adiabatic impacts (i.e., the "Dst-effect" [*McIlwain*, 1966]).

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170 3.1 SEA of MLT-dependent >100 keV electron fluxes

Figure 1 is a reexamination of the SEA undertaken by *Rodger et al.* [2016], but presented in a MLT-dependent format and constrained to -1 to +1 days around the recurrent substorm zero epoch. Figure 1 uses the trapped >100 keV electron observations from the 90e2 channel.

Many of the most interesting features seen in Figure 1 occur so quickly that they are not resolved in the -5 to +15 day format used in the earlier Rodger study, so we have limited the time axis to a much smaller range. Note that electrons drift around the Earth with increasing MLT, i.e., from top to bottom. Figure 1 includes as a white line the empirical plasmapause location, L_{pp} , determined from the AE and MLT-dependent formulation given by *O'Brien and Moldwin* [2003].

This figure demonstrates a number of strongly MLT-dependent features. From first 180 principles, one expects magnetopause shadowing to start on the dayside, and substorm 181 injections to occur on the nightside. Figure 1 is entirely consistent with those expectations. A 182 sharp decrease in outer radiation belt flux occurs when L_{pp} moves sharply inwards shortly 183 before the zero epoch time. The sudden inward L_{pp} motion will go hand in hand with the 184 inward motion of the magnetopause; both the inward motion of L_{pp} and the magnetopause 185 result from the same dynamic changes in the solar wind. When the sudden dynamic pressure 186 increase occurs in the solar wind, it brings with it an enhanced convection electric field due to 187 the increase in solar wind velocity and/or IMF intensity. This electric field enhancement will 188 result in the inward motion of the plasmapause as the balance between the magnetospheric 189 convection and corotational electric fields readjusts. The driver of the corresponding decrease 190 in flux at high L-shells is most likely a combination of magnetopause shadowing and outward 191 transport. 192

In Figure 1 the dropout is very well defined in the 15-18 MLT range, more so than 12-15 MLT (or 9-12 MLT) where it likely begins. There is also a strong suggestion in the 15-18 MLT panel that the magnetopause shadowing begins at higher L and moves inwards to lower L-shells. This is likely caused by a combination of the direct effect of magnetopause shadowing plus a cascade of losses to even lower L-shells (i.e., lower than those that are directly opened to the magnetopause after it moves inward). The lower L-shell loss cascade mentioned is likely due to very rapid outward radial transport that occurs after magnetopause shadowing; both can occur in 10's of mins to ~1 hour [*Shprits et al.*, 2006; *Turner et al.*,
2012; *Ukhorskiy et al.*, 2015].

At high *L*-shells Figure 1 shows evidence of a magnetopause shadowing induced decrease occurring at all MLT before epoch time equal zero. This is possible as the dropout drifts around the Earth faster for higher L-shells (at *L*=7 the drift period of a 150 keV electron is ~50 min). It may also reflect a delay between the solar wind driver moving L_{pp} inwards (signaling the starting of magnetopause shadowing at high *L*), and the onset of the first substorm in the cluster which defines the epoch time. Substorms can be delayed relative to the responsible change in the solar wind on global-magnetospheric convective timescales.

Substorm produced electron injections are visible in Figure 1 across a wide range of large L-209 shells on the nightside. These are seen as new sharp increases starting in the 21-24 and 0-210 3 MLT panels. It is these injections, and subsequent processes triggered by those injections, 211 which replace the electrons lost from the dropout event, and indeed, they lead to an additional 212 enhancement. We suggest that this is consistent with the importance of substorm injections, 213 rather than large scale convection, in outer belt electron acceleration, supporting the 214 conclusions of Jaynes et al. [2015] and Rodger et al. [2016]. Figure 1 demonstrates the 215 importance of considering the interaction of MLT-dependent processes when considering 216 radiation belt dynamical processes. 217

The dropouts observed by POES penetrate to low L-shells, with discernible effects down to 218 $L \sim 4$, although that feature becomes less and less obvious at those lower L-shells in the 219 statistical picture presented here. There is an ongoing debate concerning the mechanism 220 responsible for such dropouts. Some models [e.g., Ukhorskiy et al., 2015; Mann et al., 2016] 221 222 indicate that magnetopause shadowing and outward radial transport are sufficient to cause electron dropouts down to L = 4, while others [e.g., Shprits et al., 2013] conclude that some 223 additional loss mechanism, is necessary for dropouts to extend down to such low L-shells. In 224 that latter case EMIC waves causing precipitation into the atmosphere are invoked as the 225

most likely candidate. In reality, both processes are likely important and contribute differently on a case-by-case basis. Considering that perspective, it is important look for telltale signatures of both loss processes using guidelines established by previous studies [e.g., *Turner et al.*, 2014; *Xiang et al.*, 2017; *Aseev et al.*, 2017; *Shprits et al.*, 2017].

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231 3.2 SEA of MLT-dependent 52 keV proton fluxes

We follow the same processes and undertake SEA on the 52 keV trapped proton observations (from the 90P1 channel). The proton SEA allows us to test a number of predictions where behavior should be charge dependent or independent. This SEA is shown in Figure 2, in the same format as Figure 1. Protons will drift in the opposite direction to electrons, i.e., from high to low MLT, or bottom to top in Figure 2.

A number of features are shared between Figures 1 and 2. Magnetopause shadowing is 237 clearly defined for MLTs near the dayside, although we argue it evolves in the opposite way 238 to that seen in Figure 1, i.e., starting from 12-15 MLT to be very clear in the 9-12 and 6-9 239 MLT panels. This is consistent with the direction of proton drift. The magnetopause 240 shadowing starts at the time of a sharp decrease in L_{pp} , as was seen for the electrons. 241 Magnetopause shadowing is expected to be independent of particle charge, mass or energy, 242 such that electrons or protons which are drifting around the Earth on the same L-shell (but in 243 opposite directions) will encounter the magnetopause and hence be lost. The timing of the 244 dropout signature is very similar in both Figures 1 and 2, consistent with charge 245 independence and hence the magnetopause shadowing loss process - rather than, for example, 246 wave particle interactions. 247

The dynamics of protons in this energy range are less complex than those for electrons; for example, there are no significant wave-particle interactions driving acceleration, so the proton buildup is likely due to injections from the tail. This may be why the magnetopause shadowing signature is so clear and sharp in Figure 2. The figure seems to show that on the

dayside, the dropout stretches from high *L* inwards to at least L=4, where the proton flux becomes insignificant. This is clearest in the panels for 9-12 and 6-9 MLT. In contrast, for 3-6 or 0-3 MLT the dropout appears to start earlier at higher *L* and start later at lower *L*; likely due to faster drift times for higher *L*.

Figure 2 also shows increases in protons fluxes just after the zero epoch, consistent with 256 simultaneous substorm-linked enhancements of protons and electrons. These are most clear in 257 the nightside panels (03 and 21-24 MLT). The proton enhancements are clearly present down 258 to very low L-shells, at least L=2.25. We believe these are much lower L-shells than would be 259 generally expected for a substorm injection event, and note the SEA median fluxes at these L-260 shells are very low. The enhancements then progressively drift around the Earth, with the 261 injection arriving later on the morning side. As in the case for electrons, the substorm-linked 262 enhancements "refill" the magnetopause shadowing-produced dropout. One strong feature 263 present in Figure 2 is the low proton fluxes seen for all times in the 6-9 MLT region. This 264 might be indicative of the "null", or "turning point" of protons' drift trajectories. Looking at 265 the dawn quadrant outside of the Alfvén layer, protons drift trajectories will execute a sharp 266 turn around in this MLT region [e.g., Kivelson and Russell, Fig 10.25, 1995]. We suggest this 267 will result in the observed void on the dayward side of that "null" region possibly due to 268 interactions on the dayside near the magnetopause. 269

The consistency of Figure 1 and 2 in terms of timing, charge dependent drift directions, and charge independent dropouts are strong demonstrations of magnetopause shadowing produced losses, and the ability of substorms to directly compensate for those losses in these energy ranges.

4. Consequences for SPELLS events

During even mildly geomagnetically disturbed periods, electrons ranging in energy from 10's of keV up to ~1 MeV can be quickly (in a few hours or less) injected into the slot and

inner radiation belt (L<3) in events termed SPELLS [*Turner et al.*, 2017]. Evidence has been shown that the injections may serve as the dominant source of 10's to 100's of keV electrons in Earth's inner radiation belt. The physical mechanism responsible for these events, is to date, unclear.

One suggested SPELLS-production mechanism [*Lejosne et al.*, 2018] is due to subauroral polarization streams (SAPSs), which cause localized potential drops in the premidnight region. The mechanism relies on SAPS-produced electric fields helping energetic electrons to deeply penetrate the inner magnetosphere. A prediction of this mechanism is that the energetic electrons will penetrate more deeply than the energetic ions, due to their differing charge.

It has long been recognized that the POES MEPED instruments are not particularly sensitive 287 to energetic electrons [Rodger et al., 2010a; Yando et al., 2011], especially in comparison 288 with those instruments onboard the Van Allen Probes. It is clear, however, that the POES 289 data can detect the inner belt >100 keV electron enhancement which occurs during the mildly 290 geomagnetically disturbed times around recurrent substorm epochs (as seen in Figure 4 of 291 Rodger et al. [2016]). At the times of recurrent substorm epochs relatively low energy 292 (52keV) protons are injected to L-shells equivalent to the inner radiation belt, but only for 293 relatively short time periods (hours), after which their fluxes drop to the noise floor. The time 294 period of hours is consistent with losses due to charge exchange at lower L-shells. 295

One might argue that this proton injection is inconsistent with the SAPs mechanism. However, it seems important to know that the proton injection is fairly sharp and only occurs shortly after the start of the recurrent substorm epochs. In contrast, the inner radiation belt and slot region electron enhancements are not as sharply defined as the proton and outer radiation belt electron injections are. It is not totally clear that these observations rule out the SAPS argument, but neither do they support that suggestion.

We have argued that the recurrent substorm processes and SPELLS are likely to be linked. 302 Both occur during mildly disturbed geomagnetic conditions and involve electron 303 enhancements at very low L-shells. To further investigate this we undertook superposed 304 epoch analysis using the times of RBSP-observed SPELLS events as the epochs. SPELLS 305 epochs were defined as the first observation time of a sudden enhancement in electron flux in 306 the slot and inner zone by Van Allen Probes, which should bound the actual time of the 307 SPELLS to within 4.5 hours or less. There were 143 such epochs, spanning the time range 308 from 02 Dec 2012 to 20 Nov 2014. Figure 3a shows the SEA using these SPELLS epochs 309 undertaken on the daily number of SuperMAG-reported substorms [Newell and Gjerloev, 310 2011b]. The blue line shows the daily variation in daily substorm number, while the two 311 black lines show the upper and lower quartiles. During the period from the start of 2012 to 312 the end of 2014 the median number of substorms per day was 2. In contrast, on the day of the 313 SPELLS zero epoch the median substorm number is 5 (ranging from 3-7 across the quartiles). 314 This is consistent with strong substorm activity occurring in the same time period as 315 SPELLS, and potentially providing the energy for the SPELLS mechanism. 316

Finally, we show that the solar wind drivers and geomagnetic conditions around recurrent 317 substorm epochs and SPELLS epochs are very similar. We independently undertook SEA for 318 both sets of epochs, producing Figure 3b. The 4 left-hand panels are the SEA for the 143 319 SPELLS-epochs, while the 4 right-hand panels are those for the 2,052 recurrent substorm 320 epochs. In all cases 1-hour time resolution is used. From top to bottom, Figure 3 shows the 321 solar wind pressure, solar wind Epsilon parameter [Akasofu, 1981], SEA for IMF Bz, and the 322 SuperMAG-determined AU equivalent (often termed SMU). Note that the 3 solar wind 323 324 parameters were shifted to the Earth's bow shock nose.

As the left-hand panels involve ~14 times fewer epochs than the right hand panels, the SEA are much noisier than those shown for the recurrent substorm epochs. However, we argue there is close agreement in the two sets of SEA, with variations occurring on similar time scales and with very similar magnitudes. As such it seems likely that SPELLS events are
likely a subset of the recurrent SS events. It is telling that 78% of the SPELLS events can be
linked to recurrent substorms. As such, the low-*L* electron enhancements should be fairly
common during recurrent substorms, consistent with the SEA shown in Figure 4 of *Rodger et al.* [2016].

5. Summary and Discussion

In this study we have examined dynamical variations in the radiation belt during times of 334 mild geomagnetic disturbance. It is not uncommon for researchers to focus on strong storms 335 to consider changes in the radiation belts; our results demonstrate that very mild 336 disturbances are also associated with multiple processes leading to loss and enhancements 337 in flux. These changes are not limited to the outer belts, but extend even into the inner belt 338 and slot region. We have focused on time periods with high solar wind speeds where 339 clusters of substorms occur. Around these times convection from the solar wind and fresh 340 particle injections from the substorms are both present, likely leading to enhanced plasma 341 wave activity. Our results demonstrate clear signatures of magnetopause shadowing 342 impacting both electrons and protons. We have also shown how dayside magnetopause 343 shadowing losses interact with nightside substorm-produced enhancements, leading to 344 consistent MLT-dependent patterns. 345

We have shown that SPELLS events occur during the same conditions as periods of recurrent (or clustered) substorms, and that it is likely that the low-*L* electron flux enhancements are somehow linked to the substorms. Recurrent substorm epochs have been shown to display SPELLS-like enhancements. However, the enhancements are not sharply onset around the substorm cluster start, and it is certainly not clear that the SPELLS electron enhancement is directly linked to the start of the recurrent substorm cluster. Rather, these

seem to be geophysically linked. This would be consistent with SAPS as a production
 mechanism, as these are correlated with substorms.

As noted earlier, much recent progress has been provided by large, very high quality, but 354 high cost, flagship missions; such spacecraft cannot, unfortunately, provide much 355 simultaneous MLT and L coverage. However, we are now at the start of a revolution in low 356 cost access to space, alongside the development of relatively low mass, low power, but very 357 high quality radiation belt experiments. Some examples are: Colorado/LASP's CSSWE, 358 AeroCube-6, UNH/MontanaState's FIREBIRD, UCLA's ELFIN, AlbertaSat. New 359 instruments are becoming available that are suitable for deployment on cubesats - opening 360 up the possibility of large constellations to truly image the rapid dynamics of radiation belts 361 processes. Such CubeSat missions are not limited to LEO, with examples of radiation belt 362 CubeSat's planned for geostationary transfer orbit (e.g., GTOSat [Blum et al., 2018]. We 363 suggest that this study shows the value of simultaneous multi-platform observations 364 allowing MLT and L-coverage, even if with limited instrumentation. 365

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374 Data availability is described at the following websites: https://www.ngdc.noaa.gov/stp/satellite/poes/dataaccess.html (POES SEM observations). 375 http://supermag.jhuapl.edu/substorms/ (SuperMAG substorms), 376 http://supermag.jhuapl.edu/indices/ (magnetic indices and solar wind parameters from 377

378 SuperMAG). Van Allen Probes data are freely available on the science gateway site:

379 http://rbspgway.jhuapl.edu/

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514	Mark A. Clilverd, British Antarctic Survey (NERC), High Cross, Madingley Road,
515	Cambridge CB3 0ET, England, U.K. (e-mail: macl@bas.ac.uk).
516	Aaron T. Hendry, Department of Space Physics, Institute of Atmospheric Physics, Prague,
517	Czechia. (email: ath@ufa.cas.cz).
518	Craig J. Rodger, Department of Physics, University of Otago, P.O. Box 56, Dunedin, New
519	Zealand. (email: crodger@physics.otago.ac.nz).
520	Drew L. Turner, Space Sciences Department, The Aerospace Corporation, El Segundo,
521	CA, United States (email: drew.lawson.turner@gmail.com).
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Figure 1. Superposed epoch analysis of median >100 keV POES trapped electrons for the 531 Recurrent Substorm Epochs, plotted against L-shell. Each panel is for a different MLT 532 range, as labeled. Plasmapause location is shown by the white line. Note that electrons drift 533 around the Earth from top-left to bottom-right. 534

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Figure 2. As for Figure 1, but now showing the SEA for 52 keV protons. Note that protons
 drift around the Earth from bottom-right to top-left.



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Figure 3. Upper panel (a): Superposed epoch analysis of the number of SuperMAG reported substorms per day for the Van Allen Probes-observed SPELLS epochs. The blue line is the median, while the black lines are the upper and lower quartiles. Lower panels (b): Superposed epoch analysis of the number of solar wind and SuperMAG reported AU index for the Van Allen Probes-observed SPELLS epochs (lower 4 left hand panels) and the recurrent substorm epochs (lower 4 right hand panels).

Figure 1.





Time Since Epoch [days]

Figure 2.







-0.75 -0.5 -0.25 0.25 0.5 0 0.75 Time Since Epoch [days]

Figure 3.

