- 1 Electron precipitation from the outer radiation belt during the St Patrick's Day
- <sup>2</sup> storm 2015: observations, modelling, and validation

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10 **Main point # 1:** Remarkably reliable subionospheric VLF transmitter phase measurements

11 provide >30 keV electron precipitation fluxes for March 2015.

Main point # 2: VLF-inferred >30 keV electron precipitation fluxes are similar to the
 equivalent POES >30 keV loss-cone fluxes in the same region.

14 **Main point # 3:** CMIP6 >30 keV electron precipitation fluxes are only 1.3 times lower than

the VLF-inferred fluxes during the 2015 St Patrick's Day storm.

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Abstract. Recently, a model for medium energy (30–1000 keV) radiation belt-driven 17 electron precipitation (ApEEP) has been put forward for use in decadal to century-long 18 climate model runs as part of the Climate Modelling Intercomparison Project, phase 6 19 (CMIP6). The ApEEP model is based on directly observed precipitation data spanning 20 2002-2012 from the constellation of low Earth orbiting Polar Operational Environmental 21 22 Satellites (POES). Here we test the ApEEP model's ability using its magnetic local time variant, ApEEP MLT, to accurately represent electron precipitation fluxes from the 23 radiation belts during a large geomagnetic storm that occurred outside of the span of the 24

development dataset. In a study of narrow band sub-ionospheric VLF transmitter data 25 collected during March 2015, continuous phase observations have been analyzed 26 throughout the entire St. Patrick's Day geomagnetic storm period for the first time. Using 27 phase data from the UK transmitter, call-sign GVT (22.1 kHz), received in Reykjavik, 28 Iceland, electron precipitation fluxes from L=2.8-5.4 are calculated around magnetic local 29 noon (12 MLT), and magnetic midnight (00 MLT). VLF-inferred >30 keV fluxes are 30 similar to the equivalent directly-observed POES fluxes. The ApEEP\_MLT >30 keV fluxes 31 for L<5.5 describe the overall St Patrick's Day geomagnetic storm-driven flux enhancement 32 well, although they are a factor of 1.7 (1.3) lower than POES (VLF-inferred) fluxes during 33 the recovery phase.. Such close agreement in >30 keV flux levels during a large 34 geomagnetic storm, using three different techniques, indicates this flux forcing are 35 appropriate for decadal climate simulations for which the ApEEP model was created. 36

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#### 39 **1. Introduction**

One of the largest geomagnetic storms during solar cycle 24 occurred on 17 March 2015. 40 Widely known as the St Patrick's Day storm, the disturbance originated from a coronal mass 41 ejection associated with a C9 solar flare which occurred at ~02 UT on 15 March 2015. A 42 sudden storm commencement occurred at ~04 UT on 17 March due to the arrival of an 43 interplanetary shock driven by the magnetic cloud [Wu et al., 2016]. Elevated geomagnetic 44 activity levels lasted throughout 17 and 18 March, with the geomagnetic activity index Ap 45 peaking at levels of 179 nT, slowly subsiding thereafter. The storm period has been 46 investigated for many geophysical effects including the sudden loss of relativistic electrons 47 from the outer radiation belt during the early storm period [e.g., Baker et al., 2016; Shprits et 48 al., 2017]. 49

The evolution of outer radiation belt electron fluxes during large storms like the St 50 Patrick's Day storm involve a delicate balance between transport, acceleration, and loss 51 processes [e.g. Reeves et al., 2003; Glauert et al., 2018]. Gyro-resonant wave-particle 52 interactions of electrons with very low frequency (VLF) waves have been shown to produce 53 acceleration and loss within the radiation belt [Horne et al., 2016]. Waves that occur outside 54 of the plasmapause, such as VLF chorus, diffusively scatter electrons into the atmospheric 55 loss cone as well as accelerating some to higher energies [e.g. O'Brien et al., 2003]. Waves 56 inside the plasmapause, such as VLF hiss, are associated with loss processes only [e.g. 57 Meredith et al., 2006; Rodger et al., 2007]. Other waves, such as electro-magnetic ion 58 cyclotron waves have also been linked to electron precipitation over a wide range of energies 59 [e.g., Hendry et al., 2017]. 60

61 Whatever the cause of the energetic electron precipitation (EEP) into the atmosphere, the 62 generation of excess ionization at altitudes of 50-100 km affects radio communication 63 conditions, and creates odd hydrogen (HOx) and odd nitrogen (NOx) species through ion

chemistry reactions [Verronen et al., 2005]. Both HOx and NOx species are able to catalytically destroy ozone [Brasseur and Solomun, 2005], and consequently alter the radiative and dynamic balance of the atmosphere [e.g., Seppälä et al., 2009; Andersson et al., 2014, Seppälä and Clilverd, 2014]. Therefore, understanding the loss of electrons from the radiation belts during geomagnetic storms is important, not only for radiation belt dynamics, but also for understanding the effects of space weather on the climate system [Clilverd et al., 2016].

The St Patrick's Day storm has been studied previously using VLF radio signals from man-71 made transmitters. Narrowband VLF signals from naval transmitters can be received 72 subionospherically over long distances, but the phases of the received signals can vary, due to 73 a combination of changes in the transmitter-receiver path length and variations in the electron 74 density integrated along the path. Because of this, phase perturbations to quiet-day levels can 75 provide information on the characteristics of EEP into the D-region of the ionosphere 76 [Clilverd et al., 2010; Simon Wedlund et al., 2014]. Gokani et al. [2019] studied short-term 77 amplitude and phase perturbations on subionospheric paths at quasi-constant L=4 in order to 78 investigate the significance of relativistic electron precipitation into the atmosphere during 79 the first few hours of the St Patrick's day 2015 storm. The technique used in this study is 80 similar to that undertaken by Gokani et al. [2019] but here it is applied to a much longer 81 dataset, requiring high transmitter phase stability. Maurya et al. [2018] studied a 82 subionospheric path covering equatorial latitudes to show that VLF signal amplitudes were 83 perturbed for ~10 days following the storm, although analysis showing decreased D-region 84 electron densities suggested the presence of travelling ionospheric disturbances rather than 85 electron precipitation. 86

Narrow-band subionospheric VLF signals have been used to investigate the characteristics of EEP during other geomagnetic storms. Simon Wedland et al. [2014] showed that amplitude perturbations lasting 20 days occurred following a sequence of two geomagnetic

storms in July and August 2010. Enhanced outer radiation belt electron precipitation fluxes 90 over the range L = 3 to 7 with energies of 10 keV to several MeV were inferred using a 91 technique that combined the amplitude perturbations of two closely located transmitters with 92 similar frequencies. Single transmitter amplitude-only perturbations were converted to outer 93 radiation belt electron precipitation fluxes over periods of  $\sim 100$  days at a time by Clilverd et 94 al. [2010] with subsequent improvements by Neal et al. [2015]. These studies were limited to 95 ~100-day summer-only periods because of the difficulty in reproducing the observed winter-96 time quiet-time amplitude levels using modelling by the Long Wave Propagation Code 97 (LWPC) [Ferguson and Snyder, 1990]. Without knowledge of the background electron 98 density profile characteristics it is difficult to accurately model the electron precipitation 99 characteristics. Studies are preferentially limited to amplitude-only analysis much of the time, 100 due to the difficulties in determining if observed phase changes are due to geophysical, 101 transmitter, or receiver effects [Clilverd et al., 2009]. 102

Efforts to determine quiet-time D-region electron density profile characteristics over a 103 range of latitudes, including the Arctic region, have been undertaken. This is modelled 104 through the Wait profile [Wait and Spies, 1964]. Using high quality, absolute phase, multi-105 point measurements close-to, and far-from individual transmitters, the non-disturbed quiet-106 time D-region reference height (H') and sharpness (beta) parameters of the Wait profile have 107 been found for low latitudes [Thomson et al., 2014], mid-latitudes [Thomson et al., 2017] and 108 high latitudes [Thomson et al., 2018]. However, at the higher latitudes associated with the 109 magnetic field-line footprints of the outer radiation belt, this has only been achieved for 110 summer-time, daylight conditions. At the current time, the characteristics of the high latitude 111 nighttime D-region electron density that can explain observed VLF subionspheric 112 propagation signal levels remains an outstanding question. 113

A model for 30–1000 keV radiation belt driven EEP, based on satellite data, has been put forward for use in climate models [van de Kamp et al., 2016]. The EEP model is based on

electron precipitation data spanning 2002-2012 from the constellation of low Earth orbiting 116 POES satellites [Rodger et al., 2010]. The inclusion of EEP into the climate modeling inter-117 comparison project, phase 6 (CMIP6) [Matthes et al., 2017] required an EEP model that was 118 binned in geomagnetic latitude, and geomagnetic activity (the Ap index), but was zonally 119 averaged, and had a time resolution of 1 day. The model is referred to as ApEEP. Multiple 120 earlier studies into the atmospheric and climate impacts of EEP have made use of directly 121 observed POES EEP fluxes [Andersson et al., 2014; Orsolini et al., 2018; Newnham et al., 122 2018], albeit binned by time and latitude. The ApEEP model is more suitable for long climate 123 runs than the direct POES EEP flux approach [Andersson et al., 2018], as the latter is limited 124 to the time period of those direct observations. The ApEEP model incorporated in the CMIP6 125 project is suitable for climate modeling approaches back to 1850, and can be used in future 126 climate model runs, using statistically predicted Ap values [Matthes et al., 2017]. As the 127 ApEEP model is now recommended as part of the solar variability forcing set in CMIP6, it is 128 important to test the accuracy of the model output against independent datasets, as undertaken 129 in the current study. 130

Nesse Tyssøy et al. [2019] concluded that the ApEEP model >30 keV fluxes are potentially 131 too low during geomagnetic storms with Ap > 40 nT, partly because of pitch angle anisotropy 132 within the bounce loss cone (BLC). However, Rodger et al. [2013] used satellite electron 133 precipitation observations combined with ground-based riometer absorption to show that the 134 BLC was isotropic during high flux EEP events, i.e., indicating strong diffusion. Therefore, 135 there is an open question about if a large geomagnetic storm will be well represented by the 136 ApEEP model. This will depend on whether the BLC is isotropically filled by large storm-137 138 time EEP fluxes, in which case the model is likely to be correct. An updated EEP model which included 8 magnetic local time (MLT) sectors was developed by van de Kamp et al. 139 [2018], called APEEP\_MLT. An important point to note is that the ApEEP model gives the 140 same flux results as the MLT averaged ApEEP\_MLT model [van de Kamp et al., 2018]. The 141

addition of MLT sector flux information allows detailed comparison with radiation belt
 processes to be undertaken. In addition, it is now possible to make detailed comparison with
 EEP characteristics determined from ground-based subionospheric VLF narrow-band
 radiowave observations on fixed transmitter-receiver great circle paths.

In this study we analyze, for the first time, the impact of a large geomagnetic storm on the 146 phase of a transmitter continuously operating over many days. VLF transmitter phase tends to 147 be harder to measure accurately over long periods than amplitude, but it is easier to interpret. 148 High quality phase observations lasting almost a month are interpreted in terms of non-149 disturbed background ionospheric electron density profiles, and storm-induced EEP fluxes. 150 The resultant EEP fluxes are then compared with the equivalent directly observed POES 151 >30 keV loss-cone fluxes, and the output of the ApEEP\_MLT model, showing where 152 agreement exists, and where discrepancies arise. 153

## 154 **2. Geomagnetic conditions and experimental datasets**

The time variation of the geomagnetic activity index Ap for March 2015, as well as the 155 GOES-15 >800 keV and >2 MeV trapped fluxes, are shown in Figure 1. The figure shows 156 that a large geomagnetic disturbance occurred on 17 March, with Ap exceeding 150 nT for a 157 day, followed by a recovery over the next 4 to 5 days. The outer radiation belt fluxes at 158 geostationary orbit (L=6.6) show 2 to 3 orders of magnitude enhancements for both energy 159 ranges associated with the geomagnetic storm, with fluxes remaining elevated, although 160 slowly recovering, for the rest of the month (>10 days). Prior to the storm period in mid-161 March geomagnetic conditions were mostly quiet, particularly from 10 to 16 March. In that 162 time GOES-15 fluxes were slowly subsiding towards low background levels. In this study 163 data from the period 14 to 16 March are used to represent pre-storm quiet day conditions. 164

The flux of precipitating >30 keV electrons observed in the bounce-loss-cone by the POES
 SEM-2 electron telescopes [Rodger et al., 2010] are shown for the extended study period in

Figure 2. Zonal mean electron fluxes are shown for L=2 to 10, with a resolution of 0.25 L. Enhanced fluxes at L-shells less than 4 are observed following the St Patrick's Day storm on 17 March, with magnitudes reaching >10<sup>5</sup> el. cm<sup>-2</sup>s<sup>-1</sup>sr<sup>-1</sup> before slowly recovering to lower values over the next 10 days.

VLF phase data analyzed in this study was recorded by an UltraMSK receiver system [Clilverd et al., 2009] located in Reykjavik, Iceland which was set to monitor the signals from the UK Naval transmitter in Skelton (22.1 kHz, call-sign GVT). The transmitter – receiver locations are shown in Figure 3. The GVT transmitter location is indicated by the green circle, while the Reykjavik receiver is indicated by a red diamond. *L*-shell contours for L=3.5, 4, and 5.5 are shown.

The UltraMSK software uses GPS 1PPS timing to accurately determine the relative phase 177 of the GVT transmissions [Clilverd et al., 2009]. The great circle path from Skelton to 178 Reykjavik spans the *L*-shell range 2.7 to 5.4, and thus has the potential to be used to monitor 179 changes in D-region ionization conditions caused by electron precipitation from the outer 180 radiation belt. The phase analysis presented in this study is made possible because of the near 181 continuous operation of the GVT transmitter throughout March 2015, along with the 182 continuous operation of the receiver. This allows relative phase variations to be determined 183 for a period of 29 days in a row – something that is not normally possible because of 184 instability in either transmitter phase or receiver phase-lock. The transmitter amplitude was 185 also logged at Reykjavik, however the amplitude levels during geomagnetic storms, were 186 highly variable and less understandable as a monitor of long-lasting perturbations, consistent 187 with the findings of George et al. [2019] for solar flare analysis. A second receiver location is 188 also shown in Figure 3 by a red diamond close to the transmitter, at Eskdalemuir geomagnetic 189 observatory. The Eskdalemuir phase data is used to monitor the source transmitter phase prior 190 to any changes induced by ionospheric perturbations. 191

#### 192 **3. VLF phase observations**

The GVT transmitter typically goes off for a few hours of maintenance at the beginning of 193 each month, but otherwise remains on continuously, with high quality phase stability for the 194 majority of the time. The GVT relative phase variations observed from Reykjavik for March 195 2015 are shown in Figure 4. In the plot a diurnal phase variation of ~170° is apparent, 196 particularly prior to 17 March. Nighttime phase values are lower than daytime ones, with 197 rapid transitions between the two at sunrise and sunset along the great circle path between 198 transmitter and receiver. After the onset of enhanced geomagnetic activity on 17 March the 199 diurnal phase variation patterns change significantly with higher phase values both during the 200 day and the nighttime, effectively reducing the diurnal phase range to  $\sim 50^{\circ}$ . We postulate that 201 this distinct change is due to the impact of EEP on the ionosphere, affecting the 202 subionospheric VLF radio propagation. A return to more normal diurnal phase variations can 203 be observed towards the end of the month. Figure 4 also shows a representative quiet day 204 phase curve (QDC) superimposed as a red dashed line. The QDC was calculated as an 205 average of the phase on 14, 15, and 16 March. A 1º/day phase drift was applied to the QDC 206 throughout the month, as this was found to be a feature of the source transmissions, as 207 determined by the AARDDVARK receiver at Eskdalemuir, which is located close to the 208 GVT transmitter. 209

A more detailed plot of the pre-storm period on 16 March 2015 is shown in the upper panel 210 of Figure 5. GVT phase variation is given by the solid black line, while the 3-day average 211 phase variation (QDC, based on 14-16 March) is represented by the red dashed line. From 212 00-06 UT and 20-24 UT the nighttime phase values are much lower than during the daytime 213 from 08-18 UT as expected [Thomson et al., 2007]. The equivalent magnetic local time 214 215 (MLT) of the mid-point of the GVT-Iceland path is given in the upper x-axis, and we note here that there is very little difference between UT and MLT for this path (<10 minutes). The 216 super-imposed diamonds indicate the phase calculated by LWPC for the GVT-Reykjavik 217

path on 16 March, using D-region ionospheric electron number density Wait-based profiles 218 for solar zenith angle-defined beta (sharpness) and H' (reference height) values determined 219 by McCrea and Thomson [2000], and mid-latitude nighttime beta and H' values from 220 Thomson and McRae [2009]. Several features of note can be observed, including the sudden 221 phase change effects of a M2 solar flare just prior to midday (see George et al. [2019] for a 222 discussion of large solar flares and their VLF responses), and a sunrise shoulder, relative to 223 the daytime phase levels, which is caused by ozone-layer absorption of solar UV during high 224 solar zenith angle conditions [Macotela et al., 2019]. Although these two features are not 225 captured by the LWPC modelling, the close fit between the rest of the observed phase 226 variations, the QDC, and LWPC modelling results indicate a high-quality knowledge of the 227 background, undisturbed ionospheric conditions prior to the geomagnetic storm on 17 March. 228 The model-observation agreement during nighttime conditions indicates that mid-latitude 229 beta and H' nighttime values can be applied to propagation paths that do not exceed 66 in 230 latitude. This ionospheric condition knowledge provides a baseline on which to determine 231 storm-induced phase perturbation levels, and calculate the electron precipitation flux 232 involved in generating those perturbations. 233

The variation in phase during the storm onset and main phase period is shown in detail in 234 Figure 5, lower panel. The plot shows the observed phase (black line) and the QDC (red line) 235 from 16 to 21 March 2015. Shading indicates periods of nighttime on the propagation path. 236 Following the non-disturbed day on 16 March where the two lines track closely, the phase 237 shows bursts of increased phase during the daytime of 17 March, often returning to near QDC 238 levels afterwards. However, during the latter part of the day when the path experiences 239 240 nighttime conditions the phase shows a consistently large phase enhancement compared with the nighttime QDC. After 17 March the phase is continuously enhanced relative to the QDC 241 levels during the daytime and nighttime for several days, although the nighttime values can 242

be seen to be relaxing back towards the QDC from the start of 18 March. Daytime phasevalues peak on 19 March.

Average daytime (blue line) and nighttime (black line) phase perturbation levels are shown 245 in Figure 6 for the period from 9 to 31 March 2015. The phase perturbation was calculated as 246 the difference between the GVT phase and the QDC. The daytime values are averaged over 247 08-18 UT, while the nighttime values are averaged over 00-05 UT. These time ranges were 248 selected in order to minimize the impact of rapidly changing phase during sunrise and sunset 249 times, as seen in Figure 5. The nighttime phase perturbation value responds immediately 250 following the start of the storm, quickly reaching peak values of ~130°, which last for 3 251 nights before subsiding slowly towards the zero line over the next 6 nights. Daytime phase 252 perturbations increase steadily over two days, reaching a peak of ~50° before subsiding 253 slowly for the next 5 days. After the slow recovery in night- and day-time phase perturbation 254 values towards zero, from 26 March there is an additional period of elevated phase 255 perturbation levels. Any association with the S Patrick's Day storm that started on 17 March 256 is unclear. 257

**4. Modelling phase perturbations** 

With knowledge of the background D-region conditions during daytime and nighttime it is 259 possible to calculate the levels of flux of >30 keV precipitating electrons that are required to 260 generate the observed phase perturbations. Here we follow the process previously described 261 in Hardman et al. [2015], where the flux of >30 keV precipitating electrons is combined with 262 spectral gradient information via a power law scaling exponent (k) in order to generate a 263 precipitating flux from 30 to 1000 keV. A simple chemical model is then used to determine 264 the levels of excess ionization generated over a range of altitudes from 50-100 km. Finally, 265 the resultant electron number density profiles are input in to the LWPC subionospheric 266 propagation model in order to calculate the expected phase changes for a given transmitter 267

and receiver path. A full description of this process is given in Rodger et al. [2012] and Simon Wedlund et al. [2014]. As the D-region has higher electron number densities at low altitude during the daytime, compared with the nighttime, the same precipitation flux will produce different electron number density profiles at these times, and therefore different radiowave perturbation levels.

We want to invert the process described above, to calculate the flux of > 30 keV electrons 273 from the phase perturbation. In order to undertake this calculation for the St. Patrick's Day 274 storm of March 2015, we use the ambient ionospheric conditions for daytime and nighttime 275 prior to the storm obtained using the Wait profile as described in section 3, and the levels of 276 phase perturbation observed on the GVT-Reykjavik subionospheric propagation path for each 277 day and night during the storm. However, we have no ground-based experimental 278 information that would allow us to determine the spectral gradient (k), and therefore we use 279 the results from a comprehensive analysis of DEMETER electron flux observations which 280 indicate that  $k \sim -3$  for outer radiation belt fluxes during quiet geomagnetic conditions, and  $k \approx -3$ 281 3.5 for moderate/high disturbed conditions [see figure 8 in Whittaker et al., 2013]. A similar 282 power law spectral gradient analysis has also been undertaken for POES SEM-2 electron flux 283 data [van de Kamp et al., 2016; 2018] identifying similar gradient values over a wide range of 284 geomagnetic activity levels with median k ranging from -3 to -4 during nighttime, and -2 to -285 3 during the daytime, particularly for outer radiation belt fluxes where L < 5.5 as in this study 286 [see figure 3 in van de Kamp et al., 2018]. However, this study does not use the van de Kamp 287 results directly because they are already included in the ApEEP model being investigated 288 here. 289

The variation of the level of phase perturbation with imposed electron precipitation flux >30 keV is shown in Figure 7, with the upper panel representing nighttime results and the lower panel representing daytime. The nighttime panel shows phase perturbation variations for a k=-3.5 power law spectrum (solid black line) and for k=-3 and k=-4 (dashed lines). A

vertical dotted line indicates the maximum phase perturbation level, which was achieved on 294 19 March as shown in Figure 6. A horizontal red dotted line highlights the k=-3.5 flux level 295 for the 131° peak nighttime perturbation, suggesting a peak nighttime flux of  $4 \times 10^4$  el.cm<sup>-2</sup>sr<sup>-</sup> 296  $^{1}s^{-1}$  with an uncertainty of a factor of 3 above and below, introduced by the k=-3 to -4 range. 297 The panel shows that the phase perturbation levels increase smoothly with increasing flux 298 levels, and therefore during nighttime the phase perturbation is a good indicator of 299 precipitating flux levels. However, because of the combined effect of path length and electron 300 density on the received phase, the relation between flux and phase is not always necessarily 301 linear. The daytime panel shows the phase perturbation variations for a k=-2.5 power law 302 spectrum (solid blue line), with k=-2 and k=-3 results shown as dashed blue lines. The 303 vertical dotted line indicates the maximum observed phase perturbation level (52°), 304 intersecting the k=-2.5 line at the horizontal red dotted line given by a flux level of  $3 \times 10^4$  el. 305  $cm^{-2}sr^{-1}s^{-1}$ . However, the daytime phase perturbations levels show a maximum effect of ~55° 306 before reducing as higher precipitating electron fluxes are applied, leading to two possible 307 flux level results for a single phase perturbation value. This leads to a much larger uncertainty 308 in the flux, possibly as much as two orders of magnitude. 309

The daytime overturning phase issue potentially explains the relatively low perturbation 310 level determined on 18 March compared with 19 March (34° c.f. 52°, see Figure 6). Figure 7 311 suggests that instead of moderate  $\sim 10^4$  el. cm<sup>-2</sup>sr<sup>-1</sup>s<sup>-1</sup> flux levels generating the 34° daytime 312 perturbation, it could be that there are much higher fluxes involved, possibly  $\sim 10^6$  el. cm<sup>-2</sup>sr<sup>-</sup> 313  $^{1}s^{-1}$ . Clearly these overturning daytime phase perturbation levels can lead to large 314 uncertainties in any estimated flux levels for those time periods, despite the well resolved 315 316 phase changes that were observed during the storm. Error bars on VLF-phase derived fluxes shown later in the study take this uncertainty into account. 317

### 318 **5. Flux comparisons**

Having determined the response of the GVT-Reykjavik path to electron precipitation fluxes 319 we can invert this relation to convert the observed phase perturbations into an estimate of the 320 >30 keV precipitation flux during the 2015 St Patrick's Day storm period. Comparison of the 321 fluxes can be made against the electron precipitation model, ApEEP, recently published with 322 magnetic local time (MLT) included [van de Kamp et al., 2018] which can provide fluxes in 323 the region of the GVT-Reykjavik path by using the appropriate MLT zone as MLT varies 324 through the day at the longitude of the subionospheric propagation path. A further 325 comparison can be made against the POES SEM-2 fluxes [Rodger et al., 2010] measured in 326 the longitude region encompassing the GVT-Reykjavik path. 327

The upper panel of Figure 8 shows the time varying >30 keV flux determined using the 328 VLF phase measurements during the St Patrick's Day storm of March 2015. Nighttime flux 329 levels are indicated by blue asterisks, and daytime levels by blue diamonds. Vertical lines 330 indicate uncertainty ranges generated by  $\pm 0.5 k$  (see Figure 7). The directly observed POES 331 >30 keV flux levels determined from the 0° electron telescope with a 3-hourly resolution in 332 the longitude range 30° W to 15° E, averaged over the L-shell range L = 2.64 to 5.44 is 333 indicated by the black dashed line. Initially the VLF phase derived fluxes are substantially 334 lower than the POES fluxes, which is primarily due to the POES SEM-2 instrument 335 measurement noise floor of  $10^2$  el. cm<sup>-2</sup>sr<sup>-1</sup>s<sup>-1</sup> [Rodger et al., 2010], limiting the ability of the 336 satellite instrument to detect quiet or low-level precipitation fluxes. However during the 337 storm there is good agreement between the VLF phase fluxes and the POES fluxes, both 338 during the day and the night, particularly when taking into account the error bars in the VLF 339 flux. However, there are some POES flux values that are lower than the VLF-inferred fluxes, 340 341 particularly after 20th March. This may be caused by substantial flux variations occurring over small distance scales which the long-wavelength VLF technique is relatively insensitive 342 to but does influence the POES values. This would suggest that small scale precipitation 343 structure is a feature of the recovery phase of this geomagnetic storm period. 344

The lower panel of Figure 8 shows a comparison between the ApEEP MLT model output 345 (red line) and the directly observed POES >30 keV flux levels determined from the 0° 346 electron telescope with a 3-hourly resolution in the longitude range 30° W to 15° E, averaged 347 over the L-shell range L = 2.64 to 5.44 (black dashed line). Reasonable agreement between 348 the model and the POES observations occurs during the storm period, including peak flux 349 levels, and in the temporal variations throughout each day during the main phase of the 350 storm. During the recovery phase of the storm, i.e., after 20 March 2015, the ApEEP\_MLT 351 model shows a large range of flux, with very low fluxes repeating quasi-daily in the late 352 afternoon (in both UT and MLT). This is caused by much lower electron precipitation fluxes 353 occurring in the MLT afternoon sector in the observations used to build the model, 354 potentially due to the lack of whistler-mode chorus waves in this MLT sector (e.g., see Figure 355 7 of Summers et al., 1998]). . Prior to the onset of the St Patrick's Day storm the POES 356 >30 keV fluxes tend to hover around the SEM-2 instrument measurement noise floor of  $10^2$ 357 el. cm<sup>-2</sup>sr<sup>-1</sup>s<sup>-1</sup> [Rodger et al., 2010] while the ApEEP MLT model is significantly lower, 358 showing more agreement with the VLF phase results shown in the upper panel. 359

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#### 361 6. Validation

Subionospheric VLF phase measurements of the UK transmitter, GVT, made from 362 Reykjavik, Iceland, were highly reliable over almost the whole month of March 2015. The 363 reliability of the phase measurements has allowed a derivation of the energetic electron 364 precipitation fluxes generated throughout the St. Patrick's Day geomagnetic storm. Daytime 365 and nighttime electron precipitation flux derivations are made, taking into account differing 366 background D-region conditions upon which the electron precipitation generates excess 367 ionization. The electron precipitation flux (>30 keV) derived from VLF measurements can be 368 compared with equivalent directly observed POES satellite fluxes, and the ApEEP\_MLT 369 model. In all three data series the St Patrick's Day storm generated large electron 370

precipitation fluxes, with the highest levels observed during the night, and the highest
 variability observed during the day (UT and MLT).

An important question associated with the ApEEP MLT model is about its applicability 373 for use in coupled-climate model runs [Matthes et al., 2017, van de Kamp et al., 2018]. Are 374 the ApEEP predicted fluxes correct, and does the model capture the dynamics of electron 375 precipitation from large geomagnetic storms properly? The CMIP6 solar forcing dataset 376 containing the ApEEP model output provides daily average flux descriptions for input into 377 climate models [Matthes et al., 2017]. Thus in order to appropriately compare the VLF-378 derived fluxes, POES fluxes at the longitude of the Iceland-UK VLF propagation path, and 379 the ApEEP model predicted fluxes with appropriate MLT output for the same longitude 380 region, an analysis of daily average fluxes is undertaken here. 381

Nesse Tyssøy et al. [2019] concluded that the CMIP6 >30 keV fluxes are potentially under-382 represented during geomagnetic storms with large Ap, and provide a general under-estimate 383 because of the limitations of the POES electron precipitation telescope [Nesse Tyssøy et al., 384 2016]. We note that the ApEEP model predicted fluxes used in CMIP6 analysis does not 385 include any MLT variability, while the ApEEP\_MLT does. In our study the MLT version of 386 the ApEEP model is required in order to compare against the specific VLF propagation path 387 analyzed. However, van de Kamp et al. [2018] showed that the ApEEP MLT model 388 predicted fluxes, when zonally averaged, generated equivalent fluxes to the non-MLT model 389 used in CMIP6, so we can assume the conclusions of Nesse Tyssøy et al. [2019] to also be 390 valid for the ApEEP\_MLT model predicted fluxes. 391

Figure 9 shows 24-hour average >30 keV electron precipitation fluxes determined using the ApEEP\_MLT model, VLF phase perturbations, and the longitudinally restricted POES >30 keV measurements. VLF uncertainty ranges were calculated using an average of the day and nighttime uncertainty ranges shown in Figure 8. All averages are undertaken as an arithmetic mean. We note that somewhat different values could be obtained if other averaging

methods are used. The plot shows that at about the time of the geomagnetic storm (19 March 397 2015) the electron precipitation fluxes determined from the VLF phase perturbations and the 398 POES satellite show good agreement. This is consistent with strong diffusion conditions 399 filling the bounce-loss-cone isotropically [Kennel and Petschek, 1966], leading to POES 0° 400 telescope fluxes providing an accurate measurement of the precipitation flux [Rodger et al., 401 2013; Nesse Tyssøy et al., 2016]. During the recovery phase of the storm period (20-26 402 March) the POES electron precipitation fluxes show a steady decline, which is also mirrored 403 by the fluxes determined from the VLF phase perturbations. Differences between the 24 hour 404 average POES measurements and the 24 hour average VLF phase calculations during the 405 storm (i.e., from 18-25 March) are only a factor of 1.3 which might suggest that any influence 406 of non-isotropic flux distributions within the bounce-loss-cone [Rodger et al., 2013; Nesse 407 Tyssøy et al., 2016] is masked by spectral gradient uncertainties as discussed above. 408

The ApEEP\_MLT model is based on POES electron precipitation measurements organized 409 by the geomagnetic index Ap, and so some agreement is expected between the model 410 predicted fluxes and the POES longitudinally restricted measurements during this study 411 period[van de Kamp et al., 2018]. The ApEEP\_MLT model does a good job of capturing the 412 overall time variation of the storm-induced electron precipitation fluxes, and more 413 realistically exhibits a lower noise floor prior to the storm than reported by the POES fluxes. 414 The ApEEP MLT model storm-time fluxes are only about a factor of 1.7 lower than the 415 POES fluxes. Since the ApEEP\_MLT model was based on average values, it would be 416 expected to produce higher fluxes than POES for some geomagnetic storms, and lower for 417 some others. However, the close agreement in >30 keV flux levels during the St Patrick's day 418 419 storm, using the three different techniques shown here, is encouraging fory long model simulation runs (e.g., decadal climate simulations, [Matthes et al., 2017]) for which the 420 ApEEP model was created. 421

Comparison of the ApEEP\_MLT model with the fluxes determined from VLF phase 422 perturbations is made difficult because of the uncertainty in the energy spectral gradient of 423 the electron precipitation. Is the difference in flux during the St Patrick's day storm due to 424 statistical variability between the model and the VLF phase technique, or due to uncertainty 425 in the energy spectral gradient of the electron precipitation, or due to non-isotropic bounce-426 loss-cone distribution effects on POES measurements (and therefore the ApEEP\_MLT 427 model)? We could use POES measurements to calculate the energy spectral gradient of the 428 precipitating electrons, or use the information provided in the ApEEP\_MLT model [van der 429 Kamp et al., 2016; 2018]. However, there is the potential for any errors in flux level 430 determination using non-isotropic POES measurements to also influence equivalent estimates 431 of energy spectral gradient. In addition, using extra information originating from one of the 432 datasets would compromise the independence of the comparison. Thus, it is clearly more 433 reasonable to determine the energy spectral gradients independently, which we have 434 attempted here using DEMETER electron measurements. 435

436

### 437 **7. Summary**

Subionospheric VLF transmitter phase measurements have been used to infer the 438 >30 keV electron precipitation flux generated from the outer radiation belt, L<5.5, during 439 the St Patrick's Day storm of March 2015. Measurements made close to the transmitter (at 440 Eskdalemuir in Scotland) showed that the transmitted phase was constant, apart from a 1 441 degree a day systematic drift, allowing more distant observations to be used to determine 442 phase perturbations due to electron precipitation flux. Enhanced >30 keV electron 443 precipitation fluxes lasted for 8 days, with peak fluxes during the main phase of the storm 3-444 4 orders of magnitude higher than pre-storm levels, followed by a slow recovery thereafter. 445 During the extended storm period comparison between VLF-inferred >30 keV electron 446

precipitation fluxes, directly observed POES >30 keV 0° telescope fluxes, and the CMIP6
>30 keV electron precipitation flux prediction model show that:

VLF-inferred >30 keV fluxes are similar to the equivalent POES fluxes during the
 storm suggesting a weak or masked effect of non-isotropic pitch angle distributions in the
 bounce-loss-cone, particularly during high flux precipitation [Rodger et al., 2013].

2) The directly observed POES >30 keV fluxes are typically a factor of only 1.7 higher
than the CMIP6 model predicted fluxes, primarily due to higher storm-generated flux levels
during the daytime.

455 3) CMIP6 >30 keV predicted fluxes for L < 5.5 are of the same order of magnitude as the 456 VLF-inferred >30 keV fluxes in the pre-storm period, and typically 1-2 orders of magnitude 457 lower than the observed POES pre-storm fluxes.

458

The analysis presented here provides a detailed comparison between satellite >30 keV 459 electron precipitation flux measurements, VLF phase-inferred, and the CMIP6 predictive 460 flux model during one large geomagnetic storm. The finding that the CMIP6 model of 461 predicted electron precipitation (ApEEP) under-represents geomagnetic storm-time fluxes is 462 consistent with previous analysis undertaken by Nesse Tyssøy et al. [2019], although the 463 under-estimate is found to be small. Realistic electron precipitation fluxes, as inferred from 464 VLF signal analysis and POES observations during the storm, could be as large as a factor 465 of 1.7 higher than currently estimated by the ApEEP predictive model [Matthes et al., 466 2017]. The atmospheric impact of these higher flux levels in the medium energy range (i.e., 467 30-1000 keV) needs to be investigated further. Following the conclusions of this study, the 468 use of the CMIP6 model of predicted electron precipitation (ApEEP) is appropriate in terms 469 of estimating electron precipitation flux variations during geomagnetic storms. The 470 comparison done in this study used the MLT version of ApEEP rather than the zonally 471 averaged version used in the CMIP6 dataset. However, we note that the MLT version, when 472

zonally averaged, was found to be equivalent to the ApEEP version used in the CMIP6 473 dataset, but with lower quiet-time fluxes. For shorter time period runs that are made during 474 the POES observational period we recommend using EEP from the direct POES 475 measurements. Α detailed description of these datasets can be found at 476 http://chamos.fmi.fi/chamos\_apeep.html. 477

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Figure 1. A summary plot of the geomagnetic conditions and GOES-15 geostationary trapped electron flux variations (el. cm<sup>-2</sup> s<sup>-1</sup> sr<sup>-1</sup>) during the disturbed period in March 2015.



Figure 2. POES >30 keV zonal mean electron precipitation fluxes during March 2015 as a
function of *L*-shell. Data gaps are indicated by black coloring.







Figure 3. Map of the subionospheric VLF great circle propagation path from the GVT 719 transmitter in UK (green circle) to a receiver in Reykjavik, Iceland (REY, red diamond). 720 Also shown is the location of a complementary VLF receiver at Eskdalemuir in Scotland 721 (ESK) which was used to verify the transmitter stability (red diamond). Geomagnetic L-722 shell contours for are shown as solid, dashed and dotted lines. 723





Figure 4. Subionospheric VLF phase in March 2015, from the UK transmitter, GVT
(22.1 kHz), received at Reykjavik, Iceland (black line), with a superposed quiet day curve
(QDC, red dashed line) including a 1°/day phase drift caused by the transmitter.





Figure 5. Upper panel. Diurnal variations of GVT phase received at Reykjavik (black line)
on 16 March 2015, three day average QDC (red dashed line), and LWPC modelling results
(diamonds). Lower panel. The variation of GVT phase during the first few days of the St.
Patrick's Day storm (black line) compared with a QDC (red). Dark shaded times indicate

- 737 nighttime conditions, light shading indicates daytime on the VLF path. Substantial
- deviations from the QDC begin during the daytime on 17 March.



Figure 6. Average GVT phase perturbations during the nighttime and the daytime (Night
00–05 UT black line, Day 08-18 UT blue line) before, and during the St. Patrick's Day
geomagnetic storm which started on 17 March 2015.



Figure 7. Modelled GVT phase perturbation variation with >30 keV electron precipitation flux for nighttime D-region conditions (upper panel) and daytime conditions (lower panel). Solid lines represent electron energy power law spectral gradient k=-3.5, while dashed lines represent  $\Delta k$ =±0.5. Red dotted lines indicate the maximum phase perturbation level observed during the St Patrick's Day storm. See text for more details.



**Figure 8.** Upper panel. The > 30 keV flux determined using the VLF phase measurements 752 during the St Patrick's Day storm of March 2015 (blue asterisks, night; blue diamonds, 753 day). Vertical lines indicate uncertainty ranges. The black dashed line shows the POES 754 >30 keV flux levels determined from the 0° electron telescope with a 3-hourly resolution in 755 the longitude range 30°W to 15°E, averaged over the L-shell range L=2.64 to 5.44. Lower 756 panel. The ApEEP\_MLT model output (red line) for the MLT range equivalent to the 757 Scotland-Iceland VLF path and the POES >30 keV flux levels in the longitude range 30°W 758 to 15°E. See text for more details. Alternate days are shaded for clarity. 759



Figure 9. 24-hour average >30 keV electron precipitation fluxes determined using the ApEEP\_MLT predictive flux model (red line), VLF phase perturbations (blue diamonds), and the longitudinally restricted POES >30 keV measurements (black dashed line). VLF

- uncertainty ranges were calculated using an average of the day and nighttime uncertainty
- ranges shown in Figure 8.

Figure 1.



Day of month (March 2015)

Figure 2.



Figure 3.



Figure 4.



Figure 5.



Figure 6.



Figure 7.



Figure 8.





Figure 9.

