1 2	Ionospheric <i>D</i> region: VLF-measured Electron Densities compared with Rocket-Based FIRI-2018 Model
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9	Key Points:
10	• VLF-measured D region electron number densities are compared with the
11	rocket-based FIRI-2018 model over heights 60-90 km
12	• Average VLF-measured nighttime electron number densities agree very well
13	with the rocket-based FIRI-2018 model (at heights 75-90 km)
14	• By day there is reasonable agreement at low latitudes (at ~60-75 km), but at
15	high mid-latitudes FIRI lacks galactic cosmic ray effects
16	

17 Abstract

18 Ground-based very low frequency (VLF) radio propagation in the Earth-ionosphere 19 waveguide has enabled extensive electron number densities in the D region of the 20 Earth's ionosphere to be determined, by day typically below heights of 70-80 km and 21 by night in the height range ~75-90 km. Many rocket-based electron density 22 measurements have also been reported in the literature from ~ 60 km upwards using 23 current probes, and radio propagation at a few MHz between the rocket and ground. 24 Recently these rocket measurements have been summarized, and supplemented with 25 D region production-loss modeling, giving rise to a near global model named 26 FIRI-2018 (Faraday-International Reference Ionosphere) which provides electron 27 number densities as functions of height, latitude ($<60^\circ$), solar zenith angle and F10.7 28 cm solar flux. These rocket-based electron density values are here compared with 29 corresponding values from VLF measurements, by day at a low-latitude ($\sim 20^{\circ}$) and a 30 high mid-latitude (\sim 55°), and by night mainly at mid-latitudes. At night the average 31 agreement (over 75-90 km) is remarkably good. By day, at low latitude the agreement 32 is also fairly good (in the common height range $\sim 60-75$ km), with the changes with 33 solar zenith angle being moderately comparable. For daytime high mid-latitudes, the 34 agreement is less satisfactory, particularly at the lowest common altitudes, with the 35 VLF measurements showing the expected effects of cosmic rays much more than the 36 rocket-based values. Overall, we find that the *D* region description in the FIRI-2018 37 model is a significant advance on the earlier International Reference Ionosphere 38 (IRI-2016) model.

39 **1. Introduction**

40 The *D* region is the lowest part of the Earth's ionosphere, found at heights below 41 \sim 90 km. In quiet times, at non-polar latitudes, its free electrons normally extend down 42 to \sim 75 km by night and down to \sim 55 km by day. Production of electrons by day is 43 typically dominated by ionization of the minor atmospheric constituent nitric oxide 44 (NO) by direct solar Lyman- α radiation above ~65-70 km, and by cosmic galactic 45 rays ionizing all neutral constituents below this height. At night, the most significant 46 ionizing source is likely to be the indirect solar Lyman- α reradiated by the neutral 47 hydrogen in the Earth's geocorona (Banks & Kockarts, 1973). Production is 48 essentially balanced by recombination and loss processes. However, these latter 49 processes are not yet sufficiently well understood to enable calculation of reliable 50 absolute electron number densities, and so measurements are needed to create 51 empirical D region ionization models. At these low (D region) heights, satellite 52 measurements are not practical because the air density is too high creating too much drag, while the electron number densities are normally too low $(10^6 - 10^9 \text{ m}^{-3})$ to 53 54 allow sufficient reflection with a conventional ionosonde or an incoherent scatter 55 radar. Two techniques have dominated experimental electron number density 56 measurements in the D region at least at non-polar latitudes: (1) ground-based very 57 low frequency (VLF) radio propagation and (2) in situ rocket measurements. 58 59 The VLF radio propagation technique has typically used frequencies in the range 60 ~10-40 kHz radiated by powerful (and expensive) ground-based transmitters with 61 very large antennas run by the local military (to communicate with their submarines

62 near world-wide). The receivers are normally rather inexpensive with small (~1-10 m)

antennas recording amplitude and phase often continuously, and sometimes are fully

64	portable. The propagation used is in the Earth-ionosphere waveguide where the lower
65	boundary is the ground, or preferably the sea, while the upper boundary is the D
66	region of the ionosphere from which the VLF radio waves undergo partial reflection.
67	While the propagation paths over the surface of the Earth can be several thousand km
68	or more, shorter paths of a few hundred km are often more suitable for studying the D
69	region in a desired location (e.g., low latitude, mid-latitude etc.). In particular, at a
70	range of \sim 300-400 km from the transmitter, there is often a 'modal-minimum' where
71	the received amplitude has a marked minimum compared with neighboring ranges;
72	this can be thought of either as waveguide modes interfering or as due to the ground
73	wave from the transmitter interfering with (mainly) the first hop of the transmitter
74	wave reflecting from the D region (e.g., Watt, 1967).
75	
76	In this short path VLF situation, two different sub-techniques have been used. In the
77	first of these, the local geography allowed amplitude and phase to be measured as

78 functions of range over several tens of km near the modal minimum (Thomson et al., 79 2017). In the second sub-technique, the amplitude and phase were measured both near 80 (<~100 km) the transmitter (where the ground wave is dominant and the ionospheric 81 reflection is near negligible) and at ~300 km from the transmitter (near the modal 82 minimum), where there is good sensitivity to the ionospherically reflected signal (e.g., 83 Thomson et al., 2014). Both these sub-techniques avoid needing to otherwise know 84 the transmitter's radiated power and, in particular, needing to otherwise know the 85 radiated phase of the signal at the transmitter. They are both effectively measuring the 86 changes in phase and amplitude with distance which are then compared with the 87 corresponding calculated changes from the US Navy modal codes, ModeFinder 88 (Morfitt & Shellman, 1976) or LWPC (Ferguson & Snyder, 1990; see also Ferguson,

1998) using a range of possible model electron number densities versus height to find
the one that best fits the experimental observations. Both phase and amplitude
comparisons are normally essential for VLF ionospheric measurements to avoid
ambiguity and achieve good accuracy and reliability.

93

94 In contrast, the rocket measuring technique (e.g., Friedrich et al., 2018) involves 95 transmitting an MF (medium frequency) linearly polarized radio wave of a few MHz 96 usually (nowadays) from the ground for reception on the rocket where its received 97 polarization is measured as the height of the rocket changes. This is aided by the 98 rocket spinning at a slow but known rate about its vertical axis. The transmitted 99 linearly polarized wave can be thought of as being made up of two equal right and left 100 circularly polarized waves. Once these enter the plasma (i.e., where there are free 101 electrons in the ionosphere) these two travel as separate modes with different phase 102 velocities and attenuations. This means that when they 'recombine' (i.e., are measured 103 together on the rocket) their combined (quasi-linear) polarization will be found to 104 have rotated by an amount related to the amount of plasma the waves have passed 105 through (Faraday rotation) while the extent of the ellipticity of the received 106 polarization is a measure of the relative attenuation of the two polarization modes 107 (right and left) which in turn is also a measure of the amount of plasma passed 108 through. Thus the rotating polarization and the changing relative amplitudes measured 109 on the rocket are continuously measuring the plasma changes as the rocket rises (or 110 falls) and so effectively measuring the refractive index and hence electron number 111 density at the current height of the rocket.

112

113	Both the ground-based VLF-radio technique and the rocket-based MF-radio technique
114	require the (appropriate) electron-neutral collision frequency at each height to
115	determine the electron number density at that height. This requires the neutral number
116	densities at each height, particularly the nitrogen (N ₂) number density since N ₂ is
117	~78% of the neutral atmosphere at and below D region heights. Friedrich et al.
118	(2018), when reporting on their rocket-based FIRI-2018 electron density model, used
119	the NRLMSISE-00 model (Picone et al., 2002), being the latest neutral atmosphere
120	model then available. For the VLF data used here, we use the newer NRLMSIS 2.0
121	model (Emmert et al., 2020). VLF-determined D region results are normally reported
122	using the 'Wait' (e.g., Wait & Spies, 1964) height and sharpness parameters H' and β
123	(e.g., Thomson et al., 2014, 2017) to describe the variation with height of the electron
124	number density in the D region at the time and location of the VLF measurements.
125	This has been done because these two parameters can be determined largely
126	independently of the neutral density and collision frequency height profiles assumed
127	at the time. Later, as here, when the best values of electron number density at the time
128	and place of the VLF measurements are required (Thomson et al., 2018, 2021) these
129	values can be obtained retrospectively from the measured H' and β , together with the
130	most recent best estimates of the neutral atmospheric density and collision frequency.
131	
132	In determining the electron densities from the VLF observations here we use the same
133	formula for the monoenergetic electron-neutral collision frequency as used in
134	developing FIRI-2018 – i.e., $v_m = Kp$ where p is the pressure at height, h, and $K = 6.4$
135	$\times 10^5$, in SI units (Friedrich & Torkar, 1983), with $p = NkT$ being determined from the
136	appropriate NRL atmospheric model, where N is the (total) neutral number density, k
137	is Boltzmann's constant, and T is the neutral temperature. However, as discussed by

138	Thomson et al. (2018), the available VLF propagation codes, such as ModeFinder and
139	LWPC, use the Appleton-Hartree equations which assume the electron-neutral
140	collision cross section is independent of velocity whereas the rocket-based electron
141	number densities in FIRI-2018 (Friedrich et al., 2018) used the more recent Sen-
142	Wyller equations where the electron-neutral collision cross section is taken as
143	proportional to the electron velocity. This means that the above $v_m = Kp$ can be used
144	directly in the Sen-Wyller equations but requires modification before use in
145	converting the VLF-measured H' and β values into electron densities. As also
146	discussed by Thomson et al. (2018), Deeks (1966), using both Appleton-Hartree and
147	Sen-Wyller formulations at VLF, calculated a series of adjustment factors between
148	1.5 and 2.5, starting with 1.5 low in the D region and increasing monotonically to 2.5
149	high in the <i>D</i> region by which v_m should be increased to generate an effective v for
150	use in Appleton-Hartree formulations to give as near as possible the same results as
151	the Sen-Wyller formulation. Thus, as recommended by Deeks (1966), we have here
152	used this appropriate factor at each height as plotted in his figure 1a.
153	
154	As discussed in Thomson et al. (2018), Wait and Spies (1964) defined the parameter
155	$\omega_r = \omega_o^2 / v$ where ω_o is the angular (electron) plasma frequency, and v is an
156	appropriate effective collision frequency such as the "effective ν " described towards
157	the end of the previous paragraph. The electron number density is thus given by $N_e \approx$
158	$v\omega_r/3183$ (since $e^2/\varepsilon_o m_e \approx 3183$), and ω_r was taken to vary with height, <i>h</i> , as
159	$\omega_{\rm r} = 2.5 \times 10^5 \exp(h - H')\beta$ rad/s thus defining H' as the (reference) height at which
160	$\omega_r = 2.5 \times 10^5$ rad/s, and β as a (near) constant with height, but dependent, even in
161	quiet times, on latitude, time of day, and solar cycle.
162	

163	For our main comparisons here between electron number densities from FIRI-2018
164	and those from VLF measurements, we have chosen (1) daytime at a low latitude in
165	the Hawaiian Islands, latitude $\sim 20^{\circ}$ N, using transmitter NPM on 21.4 kHz, (2)
166	daytime at a high mid-latitude from the 23.4 kHz transmitter DHO in north Germany
167	(a) along the west coast of Denmark's Jutland peninsula and (b) across the North Sea
168	to Eskdalemuir, Scotland, both at latitudes \sim 55° N, and (3) nighttime at mid-latitudes,
169	all at relatively quiet times. The VLF phase and amplitude measurements were made
170	using both a (hand-held) portable loop receiver for daytime observations, and at least
171	one (usually more) separate fixed receivers recording continuously. All phases were
172	referenced to one-second pulses from GPS (the satellite Global Positioning System).
173	The fixed recorders enabled any phase or amplitude changes which occurred at the
174	transmitters to be corrected for. Details can be found in Thomson et al. (2014, 2017,
175	2018, 2021) and references therein. No polar comparisons are made because FIRI-
176	2018 does not extend higher than $\sim 60^{\circ}$ latitude. Recently Siskind et al. (2018)
177	compared some daytime VLF-derived electron number densities with those from
178	rocket soundings but did not make any detailed comparisons with the FIRI-2018
179	model.
180	
181	FIRI-2018 is available as 1980 profiles of electron number densities between 60 and
182	150 km altitude, including 11 solar zenith angles 0°, 30°, 45°, 60°, 75°, 80°, 85°, 90°,
183	95°, 100° and 130°, 5 latitudes 0°, 15°, 30°, 45° and 60°, three solar activities, $F10.7 =$
184	75, 130 and 200 sfu, and for the middle of each calendar month (i.e., a total of 11×5
185	$\times 3 \times 12 = 1980$ profiles). These model profiles are available from
186	https://figshare.com/s/357cb03b3e5bed649bbc (Friedrich et al., 2018) or
187	https://figshare.com/search?q=FIRI-2018 The profiles are actually provided from 55

188 km to 150 km, at 1 km intervals, but Friedrich et al. (2018) advise that they consider 189 them reliable only above 60 km (and for electron densities larger than 10^6 m^{-3}). 190 Recently Xu et al. (2021) have reported on parameterizing (fitting) each of these 1980 191 FIRI-2018 electron density profiles with Wait and Spies (*H'* and β) parameters. These 192 could then potentially be used to compare with the corresponding VLF-derived Wait 193 and Spies parameters. However, we prefer here to compare electron density height-194 profiles from VLF with those from FIRI-2018.

195

196 **2. Daytime Low-Latitude Comparisons: VLF and FIRI-2018**

197 The VLF results used here are those from August 2012 reported by Thomson et al.198 (2014).

199 2.1 Comparisons near Midday at Solar Zenith Angle ~10°

200 The VLF measurements made in August 2012 in Hawaii (latitude 20.5°N) described

in Thomson et al. (2014) showed that near noon when the solar zenith angle was $\sim 10^{\circ}$

202 that the D region was characterized by the Wait parameters H' = 69.3 km and $\beta = 0.49$

203 km⁻¹. These parameters lead to the electron number densities shown by the (nearly)

straight black line with large open '+' plot symbols in the upper panel of Figure 1.

205 These were calculated using the formula $N_e = v\omega_r/3183$ from section 1 above where

206 the (Deeks-adjusted) collision frequency, ν , was derived, as also explained in section

207 1, from the NRLMSIS 2.0 neutral densities for the time and location of the

208 measurements, thus representing the best estimate of the electron densities from these

209 VLF observations (i.e., using H' = 69.3 km and $\beta = 0.49$ km⁻¹).

210

211 The four curves to the left of the black VLF line are the appropriate FIRI-2018

212 (Friedrich et al., 2018) electron number density profiles for comparison, e.g., in the

213	label "M8X30L15F130", M8 means Month 8 = August, X30 means a solar zenith
214	angle of 30° ('Chi' in FIRI-2018), L15 means latitude 15°, F130 means F10.7 = 130
215	solar flux units, etc. So, three of these four curves are for solar zenith angles of 0°
216	which are likely to be the nearest match to the actual 10° , and the other at 30° is for
217	comparison. The upper orange horizontal line is at $H' = 69.3$ km, above which height
218	the VLF measurements are essentially insensitive to the electron number density (e.g.,
219	Siskind et al., 2018). This insensitivity above height, H' , applies at least as low as H'
220	≈ 58 km at the peak of an X6 flare (McRae & Thomson, 2004). The lower orange
221	horizontal line at 60 km is the height above which Friedrich et al. (2018) state that
222	their FIRI-2018 electron number densities are considered 'reliable'. As can be seen
223	the agreement between the VLF and FIRI-2018 electron number densities is good
224	near the center of the height range (64-67 km) where both sets of densities are valid,
225	but at both the lower and upper limits of the mutual height range validity (~60 km and
226	69 km) the FIRI 2018 values are lower than the VLF values by a not insignificant
227	factor of \sim 1.6. This will be further considered in section 2.3 below. In particular,
228	below heights of ~62 km, the FIRI and VLF derived electron number densities run
229	nearly parallel with a height difference of \sim 1 km meaning that the FIRI-2018 values
230	are very similar to a "Wait 70/0.5" profile ($H' = 70 \text{ km and } \beta = 0.5 \text{ km}^{-1}$); this will be
231	commented on again in the next subsection and in section 3 below.
232	2.2 Comparisons early/late in the day at Solar Zenith Angle ~75°
233	Thomson et al. (2014) showed that, when the solar zenith angle was \sim 75°, the low
234	latitude (Hawaiian) D region was characterized by the Wait parameters $H' = 78.6$ km
235	and $\beta = 0.29$ km ⁻¹ in the morning (~1730 UT) and H'=77.4 km and $\beta = 0.31$ km ⁻¹ in
236	the afternoon (~0330 UT), giving an am/pm average of $H' = 78.0$ km and $\beta = 0.30$

237 km⁻¹ for SZA=75°. These average values lead to the electron number densities shown

238	by the two (nearly) straight lines (on the left) in the middle panel of Figure 1. These
239	blue and red-brown lines with open '+' plot symbols show the electron number
240	densities using the NRLMSIS 2.0 Deeks-adjusted collision frequencies at 1730 and
241	0330 UT respectively. In this panel, the 4 curves to the right are the FIRI-2018
242	electron number densities in August at a solar zenith angle of 75° at the two latitudes
243	of 15° and 30° and at the two F10.7 values of 75 and 130 sfu (similar to the top panel).
244	Also, rather similar to the top panel, the FIRI electron number densities at the lowest
245	heights again run from $\sim 2 \times 10^6$ m ⁻³ at 55 km to $\sim 2 \times 10^7$ m ⁻³ at 62 km, i.e., are rather
246	similar to a "Wait 70/0.5" profile ($H' = 70$ km and $\beta = 0.5$ km ⁻¹); this will be
247	commented on again in section 3 below.
248	
249	The lower orange horizontal line remains at 60 km in the lower panel because this is
250	the Friedrich et al. (2018) FIRI-2018 lowest height of 'reliability'; the upper orange
251	line is now at 78 km because this is $\sim H'$ at the SZA of 75° and is thus the greatest
252	height at which the VLF measurements are sensitive to the electron number density.
253	Overall, the agreement between VLF and FIRI-2018 at SZA=75° is neither good nor
254	bad; it is, however, quite passable. More discussion on this is given in the next sub-
255	section.
256	2.3 VLF Phase & Amplitude Calculated with FIRI Compared with Observations
257	Instead of entering values of H' and β , ModeFinder allows entry of an electron density
258	versus height profile and a collision frequency versus height profile which are then
259	used to find the appropriate VLF modes under these supplied D region characteristics;
260	from these, ModeFinder calculates the amplitude and phase of the signal at the

- 261 receiver (in a very similar way to when H' and β are supplied as inputs). The FIRI-
- 262 2018 electron density profiles for August for latitudes 0°, 15° and 30° were used here.

263	A simple quadratic interpolation was used at each height (1 km intervals) to find the
264	"FIRI-2018" profile for latitude 20.5° (the latitude of the NPM-Hawaii path). This
265	was done for each of the nine (FIRI-2018 tabulated) solar zenith angles between 0°
266	and 95°, and for each of the two tabulated solar activities, $F10.7 = 75$ and 130 sfu. The
267	collision frequency profile was generated from the NRLMSIS 2.0 neutral density
268	profile, at latitude 20.5°, using the Deeks adjustment as above. The midday collision
269	frequency profile so generated was used for all 9 solar zenith angles because, at this
270	low latitude, the effect on the calculated phases and amplitudes of using more exact
271	timing was found to be negligible. The results are plotted in Figure 2 with the phases
272	shown in the upper panel and the amplitudes shown in the lower panel. The means of
273	the two F10.7 values (75 and 130 sfu) are also shown, as $F10.7 = 102$ sfu which is
274	close to the actual value of F10.7 at the time of the VLF observations (August, 2012).
275	Also shown are the best-fit curves for the VLF phases and amplitudes versus solar
276	zenith angle actually observed for both morning and afternoon from Thomson et al.
277	(2014). The FIRI-2018 model does not differentiate between morning and afternoon;
278	it specifies only the solar zenith angle.
279	

280 As can be seen, in both the phase and amplitude panels of Figure 2, there is clearly 281 some agreement between the observed values as functions of solar zenith angle and 282 those calculated from the FIRI-2018 model at least in general shape, apart from the 283 highest solar zenith angles (90°-95°). For low solar zenith angles (~10°, near midday), 284 there is a fairly significant difference with the FIRI-2018 calculated phase value of 285 ~55° being higher than that observed by ~25°. From the upper panel of Figure 2, it can 286 be seen the observed phase does not reach 55° until the solar zenith angle is 35°-40°. 287 From Thomson et al. (2014), using their figure 3 or 5, this implies the FIRI-2018

288 profile is ~0.8 km higher than the VLF profile (H' = 69.3 km) agreeing, at least

- approximately, with the apparent relative heights of the VLF and FIRI-2018 profiles
- shown in the upper panel of Figure 1 here.
- 291

3. Daytime High Mid-Latitude Comparisons: VLF and FIRI-2018

- 293 The VLF results used in this section are from July 2015 as reported by Thomson et al.
- 294 (2017), using two closely located, high mid-latitude paths from DHO in north
- 295 Germany (1) along the west coast of Jutland, Denmark and (2) across the North Sea to
- Eskdalemuir in Scotland.

297 3.1 Comparisons near Midday at Solar Zenith Angle ~33°

- 298 The VLF measurements made in July 2015 from north Germany to Denmark (latitude
- 299 ~54.5°N) of Thomson et al. (2017) showed, near noon when the solar zenith angle
- 300 was ~33°, that the D region was characterized by the Wait parameters H' = 72.8 km
- and $\beta = 0.345$ km⁻¹. These parameters lead to the electron number densities shown by
- 302 the (nearly) straight black line with large open '+' plot symbols in the upper panel of

303 Figure 3, calculated using the formula $N_e = v \omega_r/3183$ with the (Deeks-adjusted)

- 304 collision frequency, v, being derived from NRLMSIS 2.0 neutral densities for the
- time and place of the measurements. To the left, the four colored curves show for
- 306 comparison the corresponding FIRI-2018 electron number density profiles for
- latitudes near 55° N at solar zenith angles of 30° and 45°, in July 2015 when F10.7 =
- 308 130 sfu.

309

310 At the greatest heights at which the VLF measurements are sensitive (\sim 72.8 km, \sim

311 *H'*), it can be seen, in the upper panel of Figure 3, that the agreement is quite good

between the VLF-derived electron densities (at ~55° latitude) and the FIRI-2018

313	model at a solar zenith angle of \sim 33° and a latitude of 45° (FIRI data
314	"M7X30L45F130"); however, as can also be seen, the FIRI model gives somewhat
315	lower electron densities at the higher latitude of 60°. Although this overall agreement
316	at heights near 72 km is quite passable, it is also clear that towards the lowest heights
317	(~60 km) at which the FIRI model is considered reliable (Friedrich et al., 2018), its
318	electron densities are significantly lower (by a factor of nearly 3) than the VLF-
319	derived values. In particular, the FIRI electron densities at a height of 60 km are
320	nearly the same (~ 1.0×10^7 m ⁻³) at latitude 20° (as shown in Figure 1) and at latitudes
321	\sim 55° shown here, whereas the ionization due to galactic cosmic rays is expected to be
322	significantly higher at these higher latitudes (e.g., Heaps, 1978).
323	3.2 Comparisons early/late in the day at Solar Zenith Angle ~75°
324	The same set of VLF measurements reported by Thomson et al. (2017), from July
325	2015 at latitude \sim 54.5° N, in section 3.1 above, but using the path DHO to

- 326 Eskdalemuir instead of DHO to Denmark, showed that at solar zenith angles of $\sim 75^{\circ}$
- 327 the *D* region was characterized by H' = 76.6 km and $\beta = 0.27$ km⁻¹; these parameters
- 328 lead to the electron number densities shown in the lower panel of Figure 3 by the two
- 329 fairly similar nearly straight lines (to the right), with open '+' plot symbols, again
- using the formula $N_e = v \omega_r/3183$ as above. Both these lines use NRLMSIS 2.0 but the
- blue line is for 6 UT (morning) while the red-brown line is for 18 UT (afternoon) to
- 332 illustrate the small but noticeable difference. The FIRI model does not discriminate
- between morning and afternoon so such differences are not pursued further here.
- Again it needs to be noted that all the FIRI profiles in all four panels of Figures 1 and
- 335 3 are essentially the same at \sim 60 km and below, independent of latitude.
- 336

337	As can be seen in this lower panel of Figure 3, at the greatest heights at which the
338	VLF measurements are sensitive ($H' \sim 76$ km) the FIRI-2018 electron densities, for a
339	solar zenith angle of 75° are low (plot symbol squares, $\sim 1.1 \times 10^8 \text{ m}^{-3}$) by a factor of
340	\sim 2 compared with the VLF results. However, if FIRI electron densities at a solar
341	zenith angle of 60° (not shown) were used instead the agreement would be much
342	better, being then just marginally higher (~ 2.5×10^8 m ⁻³ at 76 km) than the VLF
343	results. This might possibly be due to a scarcity of rocket profiles (in solar zenith
344	angle and latitude).

345

346 Again, as in section 3.1, it is clear that towards the lowest heights (~ 60 km) at which 347 the FIRI model is considered reliable its electron densities are lower (by a factor ~ 2.5) 348 than the VLF-derived values. Also, the FIRI electron densities at a height of 60 km are nearly the same ($\sim 1.0 \times 10^7 \text{ m}^{-3}$) at latitude 20° (as shown in Figure 1) and at 349 350 latitudes ~55° shown here, whereas the ionization due to galactic cosmic rays is 351 expected to be significantly higher at these higher latitudes (e.g., Heaps, 1978). 352 Siskind et al. (2018) have also pointed out that sounding rocket profiles have tended 353 to fail to show the increasing electron densities with increasing latitude (at low 354 altitudes) expected from the well-known, corresponding increases in galactic cosmic 355 rays. 356 3.3 VLF Phase & Amplitude Calculated with FIRI Compared with Observations 357 Section 2.3 above used ModeFinder to calculate VLF phases and amplitudes on our 358 low latitude path as functions of solar zenith angle using appropriate low latitude 359 FIRI-2018 electron density profiles, and compared these with the corresponding VLF 360 observations in Figure 2. Here, in Figure 4, similar comparisons are made but for the 361 high mid-latitude path DHO to Eskdalemuir; as can be seen, the agreement between

362 the VLF observations (from Thomson et al., 2017) and the calculations using the 363 appropriate FIRI profiles for both the phases (upper panel) and the amplitudes (lower 364 panel) is relatively poor. A significant contributor to this poor level of agreement is 365 likely to be related to the FIRI-2018 model not fully recognizing a greater number of 366 galactic cosmic ray generated electrons below heights of ~ 70 km at these relatively 367 high latitudes. While these cosmic ray generated electrons will not themselves vary 368 much with solar zenith angle (during daytime), the electron density above ~ 70 km is 369 solar generated (mainly Lyman- α) and so solar zenith angle dependent. Whether this 370 effect actually shows up in the VLF phases or amplitudes depends on the exact nature 371 of the path, including its length. Here, as the VLF observations show, the VLF phase 372 is quite strongly solar zenith angle dependent while the amplitude is not.

373

4. High Mid-Latitude VLF Electron Densities below ~60 km Height

375 Comparing the two panels of Figure 3, it can be seen that at heights below ~ 60 km the 376 VLF-derived electron number densities near noon are lower than early or late in the 377 day; e.g., at 55 km specifically, the VLF-derived electron number density at a solar zenith angle of 75° shows $\sim 1.2 \times 10^7$ m⁻³ in the lower panel while at a solar zenith 378 angle of ~33°, in the upper panel, it shows as ~ 0.9×10^7 m⁻³, i.e., ~30% lower. This 379 380 seems unlikely because galactic cosmic rays, which are the principal ionizing source 381 at these heights, are not solar zenith angle dependent. Another possibility could be the 382 photo-detachment, by visible light at sunrise, of electrons (e.g. Kazil et al., 2003, 383 Ogawa & Shimazaki, 1975) from negative ions generated during the night from 384 galactic cosmic rays, but this effect is likely to have faded when the sun has risen to a 385 solar zenith of 75° (and also does not occur at dusk).

386

387	To assess whether this difference is significant in terms of likely errors, the noon
388	profile of $H' = 72.8$ km and $\beta = 0.345$ km ⁻¹ as determined by Thomson et al. (2017)
389	was compared with a similar profile with two parts, $H' = 76.6$ km and $\beta = 0.27$ km ⁻¹
390	(as for a solar zenith angle of 75°) below 60 km and $H' = 72.6$ km and $\beta = 0.355$ km ⁻¹
391	above 60 km as shown in the upper panel of Figure 5. The lower two panels of Figure
392	5 here are taken from Thomson et al. (2017) showing the measured (data points) and
393	calculated (lines) phases and amplitudes versus distance, with the new (ModeFinder)
394	calculated (thick green) lines superposed for the new two-part profile with $H' = 72.6$
395	km and $\beta = 0.35$ km ⁻¹ (for phase) and $\beta = 0.36$ km ⁻¹ (for amplitude) above 60 km and
396	$H' = 76.6$ km and $\beta = 0.27$ km ⁻¹ below 60 km. As can be seen the (previous) best fit
397	black lines ($H' = 72.8$ km and $\beta = 0.345$ km ⁻¹) fit the data points just as well as the
398	new green lines. The change of just 0.2 km in H' and 0.01 km ⁻¹ in β is within the
399	originally estimated experimental error and so is not of much significance.

400

401 Of course, instead of adjusting the noon profile below 60 km to the 75° profile, the 402 75° profile could have been adjusted instead (e.g., by similarly making into two parts) so that, below 60 km it matched the original, H' = 72.8 km and $\beta = 0.345$ km⁻¹, noon 403 404 profile, or indeed both noon and 75° profiles could have each been even more slightly 405 adjusted (in two parts) so that they matched below 60 km. This would involve < 0.2km in H' and <0.01 km⁻¹ in β and so be of near negligible significance. Also new in 406 407 Figure 5 are orange '+' symbols, identical to the new green profile but with it cut-off 408 at a height just below 57 km, showing that the VLF technique is here sensitive down 409 to heights at least as low as 57 km.

410

411

412 5. Nighttime *D* region Comparisons: VLF and FIRI-2018

413 The nighttime VLF-derived electron number densities used in this section come from 414 long, nearly all-sea paths over a range of latitudes, but excluding both the polar regions (>~60° latitude, as does FIRI-2018) and the equatorial regions (say <~15° 415 416 geomagnetic latitude), mainly from Thomson et al. (2007) but also supported by 417 Thomson and McRae (2009). Thomson et al. (2007) found $H' = 85.1 \pm 0.4$ km and β $= 0.63 \pm 0.04$ km⁻¹ on average, predominantly in summer and under conditions nearer 418 419 solar minimum than solar maximum. The FIRI-2018 profiles compared in this section 420 are those with solar zenith angles of 130° (the highest solar zenith angle reported in 421 FIRI-2018) which correspond with being full night. The next highest solar zenith 422 angle available in FIRI-2018 was 100°; it was decided not to use these 100° values for 423 comparisons because they would be only just into full darkness and the VLF data 424 reported in Thomson et al. (2007) was wholly or mainly with solar zenith angles 425 greater than 100°, i.e., with the ionosphere well settled into full night. 426 427 Figure 6 compares the VLF-derived nighttime electron number densities with those 428 from FIRI-2018. The thick black solid (nearly) straight line with the open '+' plot 429 symbols shows the VLF-derived electron number densities for H' = 85.0 km and $\beta =$ 430 0.65 km⁻¹ using collision frequencies derived (as above) from the NRLMSIS 2.0 431 model (nominally for July at 45° latitude, 0° longitude, and 0 UT, though this was not 432 critical). The solid blue line is similarly calculated for H' = 85.0 km and $\beta = 0.60$ 433 km⁻¹. The FIRI-2018 electron number densities are shown with the dot/dashed curves 434 with the plot-symbol labeling similar to that used above except that the 'X130' 435 (indicating a solar zenith angle of 130°) is not shown in the labels because they are all 436 for 130°. So, e.g., 'M7L45F130' means Month 7 (i.e., July), Latitude 45°, solar Flux

437	130 sfu. The VLF data correspond to solar activity roughly midway between F10.7
438	solar fluxes of 75-130 sfu. Friedrich et al. (2018) have indicated that FIRI-2018
439	electron number densities below 1×10^6 m ⁻³ are not considered reliable; so these have
440	been included only for completeness. At night (when $H' \approx 85$ km), ModeFinder
441	calculations show that the VLF technique is sensitive down to heights \sim 75 km, and up
442	to heights of ~90 km, i.e., somewhat above H' (the daytime sensitivity upper height
443	limit mentioned in section 2.1).
444	

445 In Figure 6, it can be seen that the agreement between the night VLF-derived 446 D region electron number densities and those from the (rocket-based) FIRI-2018 447 model is remarkably good, especially at latitudes near 45°. It must be emphasized that 448 these are all averaged values; the night D region is rather variable (e.g., Thomson et 449 al., 2007) in time and space (compared with the daytime D region) whether observed 450 by ground-based VLF or by rocket-borne MF radio. In particular, as pointed out by 451 Friedrich et al. (2018), the actual individual nighttime electron density profiles as 452 measured by rockets at a particular place and time often show 'ledges' where the 453 electron density increases much more rapidly with height than shown here in Figure 454 6, but these ledges occur (randomly) at different heights, typically between ~80-90 455 km, and so, when averaged, result in the profiles in Figure 6. At night, galactic cosmic 456 rays generate free electrons at similar rates as by day; these electrons are then 457 similarly rapidly removed by attachment to neutral molecular oxygen molecules: e⁺ + $O_2 \rightarrow O_2$. By day, photons of (visible) sunlight immediately release these electrons 458 459 again from the negative ions, resulting in cosmic ray generated free electrons 460 dominating below heights of 65-70 km as noted above. In darkness, this last step

461 cannot occur, resulting in very few free electrons from cosmic rays at night (Thomson462 et al., 2021; see also Banks & Kockarts, 1973).

463

Cummer et al. (1998) and Cheng et al. (2006) used VLF propagation from natural

lightning (sferics) to determine nighttime D region characteristics over the USA at

latitudes ~35°-37° with both also finding $H' \approx 85.0$ km but both finding significantly

467 lower values of β with Cummer et al. finding $\beta \approx 0.50$ km⁻¹ and Cheng et al. $\beta \approx 0.45$ 468 km⁻¹.

469

470 **6.** Neutral Atmosphere Collision Effects on VLF Electron Densities

471 The VLF-derived electron number densities presented in the previous sections here

472 were calculated from their measured H' and β values using the NRLMSIS 2.0 neutral

473 density model (Emmert et al., 2020) as explained in section 1 above. These VLF-

474 derived electron number densities are now compared with (1) those similarly derived

475 from the NRLMSISE-00 model (Picone et al., 2002) as used by Friedrich et al. (2018)

476 for FIRI-2018, and (2) those derived using, as also discussed in section 1, the 'Wait'

477 collision frequency (with collision cross section independent of velocity, and so with

478 no Deeks adjustments) as incorporated in ModeFinder, i.e., $v = 5.0 \exp(-0.15(h - 70))$

479 MHz. In Figure 7, the top and middle panels make these comparisons for the low

480 latitude case (~20° N) and the high mid-latitude case (~55° N) respectively, at the two

481 appropriate daytime solar zenith angles (mid-day, 10° and 33° respectively, and 75°),

482 while the bottom panel is for night (at mid-latitudes).

483

484 It can be seen that the electron densities using NRLMSISE-00 are normally only very

485 slightly (and so negligibly) larger than those using NRLMSIS 2.0. Only at both higher

486	latitude (\sim 55° N) and higher solar zenith angle (SZA=75°), and at daytime heights >70
487	km, do the NRLMSISE-00 electron densities become appreciably, though still
488	marginally, larger (by $\sim 10\%$) than those using NRLMSIS 2.0. In contrast, while the
489	'Wait'-derived electron densities nearly always deviate from the NRLMSIS 2.0
490	values more than those from NRLMSISE-00, none-the-less they do not deviate very
491	much for both daytime low latitudes (~20° N) and nighttime mid-latitudes. However,
492	as can be seen in the middle panel, the 'Wait'-derived electron densities at the high
493	mid-latitude of ~55° are typically significantly lower by up to factors of ~1.6,
494	compared with those derived from the more modern collision frequencies and
495	atmospheric models.

496

497 7. Comparisons: IRI and VLF

498 The D region part of the FIRI-2018 model is considered a significant advance on the 499 D region part of the current (2016) International Reference Ionosphere (IRI-2016) 500 model (Bilitza, 2017). Here we compare the VLF electron number densities with the 501 equivalent IRI-2016 profiles. In the D region, IRI-2016 uses IRI-95 which is based on 502 an older and much smaller selection of typical rocket profiles such as those in 503 Mechtly et al. (1972) who also used Langmuir probes calibrated by Faraday and 504 differential-absorption data. The IRI-2016 data used here came from the NASA web 505 page https://ccmc.gsfc.nasa.gov/modelweb/models/iri2016 vitmo.php which also 506 includes a second D region model "FT-2001" which is a much earlier version of FIRI 507 (Friedrich & Torkar, 2001). Figure 8 shows comparisons between (1) the VLF-508 derived electron number densities presented here in sections 2, 3 and 5, and (2) the 509 corresponding IRI data. The top and middle panels are for mid-day at latitudes of 510 $\sim 20^{\circ}$ and $\sim 55^{\circ}$ respectively, while the bottom panel is for night at mid-latitudes. As

511	can be seen, a principal disadvantage of the IRI-2016 D region models is that their
512	electron density profiles do not extend down as low as either the FIRI-2018 or the
513	VLF profiles. By day (upper two panels) the IRI-2016 profiles extend down to ~ 65
514	km compared with down to ~ 60 km for FIRI-2018. This means they are essentially
515	not registering electrons generated by galactic cosmic rays; it also largely precludes
516	their use for VLF modeling such as for the phase and amplitude calculations
517	performed using FIRI-2018 here in sections 2.3 and 3.3. Similarly at night (bottom
518	panel) it can be seen that the IRI-95 and FT-2001 profiles extend down only to 80 and
519	83 km respectively. Also, although the agreement between the VLF and FT-2001
520	profiles is good (down to just ~83 km), there is little agreement with IRI-95 profile. In
521	the top panel, even where the red, IRI-95 profile crosses the black VLF line, the IRI-
522	95 line has a slope corresponding to $\beta = -0.40 \text{ km}^{-1}$ while the VLF observations give
523	$\beta = -0.49 \text{ km}^{-1}$ which, as can be seen in figure 3 of Thomson et al. (2014), is quite a
524	large difference (corresponding to 4-5 standard deviations or 4-5 dB in amplitude).
525	

526 8. Summary and Conclusions

527 The FIRI-2018 model uses rocket observations and modeling to provide an 528 impressive range of ionospheric electron number densities for heights, in 1-km 529 intervals, from 150 km down to nominally 55 km though only those above 60 km, and $>10^6$ m⁻³, are considered reliable (Friedrich et al., 2018). Such profiles are provided at 530 531 5 latitudes (0-60°), 11 solar zenith angles, 3 solar activities and for each of 12 months, 532 making a total of 1980 profiles. While a significant amount VLF data exists, this has 533 not yet been sorted comprehensively. We therefore chose to select three 534 representative conditions where extensive quality VLF data were available: nighttime 535 at mid-latitudes, daytime at low latitude ($\sim 20^\circ$) and daytime towards the higher end of

536	the FIRI-2018 latitudes (~55°). While VLF data can give very good electron number
537	densities, the technique works only in the lowest parts of the Earth's ionosphere.
538	During quiet times this means \sim 75-90 km altitude by night, \sim 55-70 km near mid-day
539	at low latitudes, and \sim 60-75 km at higher solar zenith angles or mid-latitudes. Both
540	the VLF technique and this lowest ionospheric region are none-the-less very
541	important in geophysics. Both man-made and natural (e.g., lightning) VLF radio
542	signals propagate in the Earth-ionosphere waveguide and so are sensitive to the lower
543	D region, by day, night, dawn and dusk, both in quiet conditions and when disturbed,
544	such as by energetic electron (or proton) precipitation or by solar flares or by (extra-
545	terrestrial) gamma rays etc. Production and loss modeling in the D region is very
546	important in attempting to get a good quantitative understanding of the key
547	mechanisms there. Recently Siskind et al. (2018) found that, by day, their theoretical
548	photochemical model agreed better with VLF-derived electron densities below 68-70
549	km than the corresponding electron densities from rockets. There appear to be
550	multiple complex processes involved in such D region modeling which are likely to
551	result in further developments but need the support of measured electron number
552	densities.
553	
554	The agreement found here between VLF-derived and FIRI-2018 electron densities is,
555	perhaps surprisingly, best at night. This may be partly because, at night, both

techniques have good sensitivity in the height range 75-90 km. The MF rocket

sensitivity has a threshold of typically between 10^9 m^{-3} at 60 km and 10^7 m^{-3} at 80 km

558 (Friedrich et al., 2018; Jacobsen & Friedrich, 1979) and so in the relevant height

range of 75-90 km the more sensitive but much less certain Langmuir probe data will

560 have been only minimally required, if at all. Also, the lower limit for electron

densities of 10⁶ m⁻³ (at any height) will have been encountered only below ~75 km. In
addition, both techniques will have averaged out any of the "ledges" mentioned here
in section 5: VLF by averaging over long paths and by averaging over several days,
and FIRI-2018 by averaging over a number of rocket profiles.

565

For daytime low latitudes ($\sim 20^{\circ}$), the agreement found here between VLF-derived and FIRI-2018 is quite reasonable, specifically for the electron densities at the solar zenith angles of $\sim 10^{\circ}$ and $\sim 75^{\circ}$ and more generally in terms of the agreement between observed and ModeFinder+FIRI-2018 predicted VLF phases and amplitudes as

570 functions of solar zenith angle. It might well be fair to say that probably from the very

571 general, global perspective of the FIRI-2018 model, the agreement is quite good,

572 while from the highly focused VLF perspective, the agreement is more modest.

573

574 For daytime high mid-latitudes (\sim 55°), the degree of agreement found here between 575 VLF-derived and FIRI-2018 is at best only quite modest, with the FIRI-2018 densities 576 being consistently lower by up to a factor of \sim 3 than the VLF-derived values. Again, 577 it might well be fair to say that probably from the very general, global perspective of 578 the FIRI-2018 model, the agreement is generally passable while, from the highly 579 focused VLF perspective, the agreement is fairly marginal. A significant factor, at 580 these higher latitudes in daytime, may well be that FIRI-2018, at the lowest altitudes, 581 in compromising between the rocket (probe, rather than MF) observations and 582 modeling, has perhaps underestimated the galactic cosmic ray generated electrons. In 583 contrast, the FIRI-2018 profiles at low latitude (20°) and high solar zenith angle (75°) , 584 in the middle panel of Figure 1, show an electron density 'bulge' at ~65 km, which 585 appears consistent with cosmic ray generated electrons, below a 'dip' at 70 km down

to which height Lyman-α is likely no longer penetrating at this high solar zenithangle.

588

589	Although, as mentioned above, the FIRI-2018 model does not extend above latitudes
590	\sim 60°, comparisons of VLF-derived electron densities in polar regions with those from
591	sounding rocket profiles and an MF radar have shown them to be quite comparable
592	both by day and by night (Thomson et al. 2018, 2021).
593	
594	Overall, our comparisons here with specific VLF-derived electron number densities
595	indicate that the FIRI-2018 model, though under-estimating at higher latitudes by day,
596	is a major advance at least compared with IRI-2016. We conclude that FIRI-2018
597	reasonably represents the D region electron number density profiles on a near global
598	basis (i.e., at latitudes below 60°), at least in the low height ranges where it can be
599	compared with VLF measurements.

600 Data Availability Statement

- 601 The NRLMSIS 2.0 neutral atmospheric data were obtained using the FORTRAN code
- 602 at <u>https://map.nrl.navy.mil/map/pub/nrl/NRLMSIS/NRLMSIS2.0/</u> but should also be
- 603 available from https://kauai.ccmc.gsfc.nasa.gov/instantrun/msis
- 604 The NRLMSISE-00 neutral atmospheric data were obtained from the NRLMSISE-00
- 605 web model at <u>https://ccmc.gsfc.nasa.gov/modelweb/models/nrlmsise00.php</u> but
- 606 should also be available from https://kauai.ccmc.gsfc.nasa.gov/instantrun/msis
- 607 The FIRI-2018 model profiles are available from
- 608 https://figshare.com/s/357cb03b3e5bed649bbc (Friedrich et al., 2018) or
- 609 <u>https://figshare.com/search?q=FIRI-2018</u>
- 610 The VLF data used here came from Thomson et al. (2007, 2014, 2017).
- 611 US Navy code LWPC is available at <u>https://github.com/mlhutchins/LWPC</u>
- 612 The US Navy computer program referred to here as ModeFinder is a slightly modified
- 613 version of MODEFNDR (e.g., Thomson, 1993; Nunn & Strangeways, 2000) and
- 614 MODESRCH described and listed in Morfitt & Shellman (1976).
- 615 Solar zenith angles were determined at <u>https://gml.noaa.gov/grad/solcalc/</u>
- 616 The IRI data used here came from the NASA web page:
- 617 https://ccmc.gsfc.nasa.gov/modelweb/models/iri2016 vitmo.php
- 618

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- 622 Weather Observatory.
- 623

624 **Figure Captions**

626 Figure 1. Low latitude (~20° N) comparisons of VLF-derived (using 21.4 kHz) and

- 627 FIRI-2018 electron number densities as functions of height. (Top panel) At solar
- 628 zenith angles $\sim 10^{\circ}$. (Middle panel) At solar zenith angles $\sim 75^{\circ}$. (Bottom panel)
- 629 Illustrates FIRI-2018 plot symbol labels such as "M8X75L30F130". VLF plot symbol
- 630 labels such as "69.3/0.49 NRL2.0" mean H' = 69.3 km and $\beta = 0.49$ km⁻¹, and NRL2.0
- 631 means the NRLMSIS 2.0 neutral atmosphere model was used to determine the
- 632 collision frequencies.
- 633

625

Figure 2. VLF phases and amplitudes, as functions of solar zenith angle, observed

and calculated for the short (306-km), nearly all-sea, low latitude path from NPM

636 (21.4 kHz) in Hawaii. The thick dashed lines are the observations (from Thomson et

al., 2014), light blue for morning and dark red for afternoon. The green, black and red

data points, joined by lines of the same colors, were calculated with ModeFinder

- 639 using the corresponding FIRI-2018 profiles (interpolated for latitude and F10.7 as
- 640 appropriate see text for details).
- 641

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642 Figure 3. High mid-latitude (~55° N) comparisons of VLF-derived (using 23.4 kHz)
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and FIRI-2018 electron number densities as functions of height. Plot symbol labels

are explained in Figure 1 above. (Top panel) Comparisons at solar zenith angles ~33°.

645 (Bottom panel) Comparisons at solar zenith angles ~75°.

646

Figure 4. VLF phases and amplitudes, as functions of solar zenith, observed and
calculated for the 748-km nearly all-sea, high mid-latitude path, from DHO (23.4)

649 kHz) to Eskdalemuir. The observations (from Thomson et al., 2017) are shown in

points, joined by lines of the same colors, were calculated with ModeFinder using thecorresponding FIRI-2018 profiles.

653



- 676 panel) Daytime at low latitudes (~20° N). (Middle panel) Daytime at high mid-
- 677 latitudes (~55° N). (Bottom panel) Nighttime at mid-latitudes.

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Figure 1.



Figure 2.



Figure 3.



Figure 4.



Figure 5.



Figure 6.



Figure 7.



Figure 8.

