

Daytime D-region Parameters from Long Path VLF Phase and Amplitude

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13 **Abstract**

14 Observed phases and amplitudes of VLF radio signals propagating on very long paths are used to
15 validate electron density parameters for the lowest edge of the (D-region of the) Earth's ionosphere at
16 low and midlatitudes near solar minimum. The phases, relative to GPS 1-sec pulses, and the
17 amplitudes were measured near the transmitters (~100-150 km away), where the direct ground wave
18 is dominant, and also at distances of ~8-14 Mm away over mainly all-sea paths. Four paths were
19 used: NWC (19.8 kHz, North West Cape, Australia) to Seattle (~14 Mm) and Hawaii (~10 Mm),
20 NPM (21.4 kHz, Hawaii) and NLK (24.8 kHz, Seattle) to Dunedin, N.Z. (~8 Mm and ~12 Mm). The
21 characteristics of the bottom edge of the day-time ionosphere on these long paths were found to
22 confirm and contextualize recently measured short-path values of Wait's traditional height and
23 sharpness parameters, H' and β , after adjusting appropriately for the (small) variations of H' and β
24 along the paths due to (1) changing solar zenith angles, (2) increasing cosmic ray fluxes with latitude,
25 and (3) latitudinal and seasonal changes in neutral atmospheric densities from the (NASA) MSIS-E-
26 90 neutral atmosphere model. The sensitivity of this long-path (and hence near global) phase-and-
27 amplitude technique is $\sim\pm 0.3$ km for H' and $\sim\pm 0.01$ km⁻¹ for β , thus creating the possibility of treating
28 the height ($H' \sim 70$ km), as a fiduciary mark (for a specified neutral density) in the Earth's atmosphere
29 for monitoring integrated long term (climate) changes below ~70 km altitude.

30

31 **1. Introduction**

32 The lowest altitude part of the Earth's ionosphere is the D-region. In this region the neutral
33 atmosphere is ionized mainly by solar EUV radiation and galactic cosmic rays. Low in the D-region,
34 the down-going solar EUV radiation is increasingly absorbed by the increasing atmospheric density;
35 also the electron attachment and recombination rates become so high that the free electron density
36 becomes very small. The lower D-region (~50-75 km) forms the rather stable upper boundary, or
37 ceiling, of the Earth-ionosphere waveguide while the oceans and the ground form the lower boundary.
38 Very Low Frequency (VLF) radio waves (~3-30 kHz) travel over the Earth's surface in this
39 waveguide. Observations of the propagation parameters of these waves result in one of the best
40 probes available for characterizing the height and sharpness of the lower D-region. The (partial)
41 ionospheric reflections of the VLF waves occur because the electron densities (and hence refractive
42 indices) change rapidly (in the space of a wavelength) with height in this region (~50-75 km)
43 typically from less than $\sim 1 \text{ cm}^{-3}$ up to $\sim 1000 \text{ cm}^{-3}$, near midday. These electron densities are not
44 readily measured by means other than VLF. Reflected amplitudes of higher frequency radio signals,
45 such as those used in incoherent scatter radars, tend to be too small and so are masked by noise or
46 interference. **The air density at these heights is too high for satellites, causing too much drag,**
47 **but too low for balloons, providing too little buoyancy.** Rockets are expensive and transient;
48 although some have given good results, there have generally been too few to cope with diurnal,
49 seasonal and latitudinal variations.

50 Because VLF radio waves penetrate some distance into seawater and, because they can be readily
51 detected after propagating for many thousands of km, the world's great naval powers maintain a
52 number of powerful transmitters to communicate with their submarines. The phase and amplitude of
53 the received signals provides a good measure of the height and sharpness of the lower edge of the D-
54 region. The US Naval Ocean Systems Center (NOSC), developed the two computer programs,
55 **'ModeFinder' (also known as 'MODESRCH' or 'MODEFNDR'), and 'LWPC' ('Long Wave**

56 **Propagation Capability')** which take the input path parameters, calculate appropriate full-wave
 57 reflection coefficients for the waveguide boundaries, and **search for those modal angles which give**
 58 **phase changes of integer multiples of 2π across a full traverse of the guide (both up and down,**
 59 **after reflection from both upper and lower boundaries)**, taking into account the curvature of the
 60 Earth [e.g. *Morfitt and Shellman*, 1976; *Ferguson and Snyder*, 1990]. Further discussions of the
 61 NOSC waveguide programs and comparisons with experimental data by the US Navy and others can
 62 be found in *Thomson* [1993, 2010], *McRae and Thomson* [2000, 2004], and references therein.

63 The NOSC programs can take arbitrary electron density versus height profiles supplied by the
 64 user to describe the D-region profile and thus the ceiling of the waveguide. However, from the point
 65 of view of accurately predicting (or explaining) VLF propagation parameters, this approach
 66 effectively involves too many variables to be manageable in our present state of knowledge of the D-
 67 region. As previously, we follow the work of the NOSC group by characterizing the D-region with a
 68 "Wait ionosphere" defined by just two parameters, the 'reflection height', H' , in km, and the
 69 exponential sharpness factor, β , in km^{-1} [*Wait and Spies*, 1964]; the studies referenced in the previous
 70 paragraph also found this to be a satisfactory simplification.

71 Daytime propagation is rather stable, potentially resulting in well-defined values of H' and β
 72 characterizing the lower D-region. ModeFinder and LWPC allow users to supply appropriate values
 73 of H' and β to determine the amplitude and phase changes along the path and so compare with
 74 observations. For the short (~ 300 km) low-latitude path, from NWC to Karratha, on the coast of N.W.
 75 Australia ($\sim 20^\circ\text{S}$ geographic, $\sim 30^\circ\text{S}$ geomagnetic, see Figure 1), *Thomson* [2010] used VLF
 76 observations plus ModeFinder to determine $H' = 70.5$ km and $\beta = 0.47$ km^{-1} near midday in late
 77 October 2009 (i.e. with the **Sun** near the zenith). Similarly, for the short (~ 360 km) high-midlatitude
 78 path, NAA (Maine, USA) to Prince Edward Island, Canada ($\sim 46^\circ\text{N}$ geographic, $\sim 53.5^\circ\text{N}$
 79 geomagnetic), *Thomson et al.* [2011] used VLF observations plus ModeFinder to determine $H' = 71.8$
 80 km and $\beta = 0.34$ km^{-1} near midday in June/July 2010 (i.e. with the **Sun** again near the zenith). The

81 lower β at the higher latitude site was attributed to the much higher galactic cosmic ray fluxes at
82 higher latitudes and **enabled a tentative plot of β** versus geomagnetic latitude to be produced.

83 In the current study here, we use phase and amplitude changes observed along very long near all-
84 sea paths to check on and, to some extent improve on, these values of H' and β . The short paths were
85 needed to measure variations (particularly in β) with latitude. However, although considerable effort
86 was used to try to have these short paths as near all-sea as possible (and hence avoid the considerable
87 uncertainties of land, particularly its low conductivity), the reality is that all the available transmitters
88 are on land. Receiving is also done much more conveniently on land. For modeling purposes, both the
89 low-latitude short path and the high-midlatitude short path were treated essentially as all-sea on the
90 assumption that the parts of the paths that were over land were close (~ 10 km) to the sea and so likely
91 to have near sea-conductivities. The use of long, nearly all-sea paths used here enables this previous
92 nearly all-sea assumption for the short paths to be checked and validated, because the proportion of
93 the path over land on the long paths here is not only much lower but also the bulk of the paths are far
94 from land (unlike the short paths which tend to pass along and close to coastlines even when over the
95 sea).

96 Of course, a disadvantage of long paths (in contrast to short paths) is that allowance needs to be
97 made for changes in some of the waveguide parameters along the length of the path. LWPC and
98 ModeFinder generally give very similar results but, because LWPC is set up to automatically take
99 into account changes in the geomagnetic dip and azimuth along the path, it is used for the long paths
100 here. Changes in H' and β due to changing solar zenith angle along the path can be found from
101 *Thomson* [1993] and *McRae and Thomson* [2000], while changes in β due to changing geomagnetic
102 latitude can now also be allowed for from the plot in *Thomson et al.* [2011] mentioned above.
103 Changes in H' with latitude and season depend effectively on the height changes of a fixed neutral
104 density near 70 km altitude and can be estimated from the MSIS-E-90 neutral atmospheric density
105 model [http://omniweb.gsfc.nasa.gov/vitmo/msis_vitmo.html]. Thus it is only now that we are able

106 to make a detailed study of long paths where propagation conditions vary significantly with distance
107 along the path. A clear advantage of long paths (in addition to being able to have a very low
108 proportion of land) is that not only are there much greater phase and amplitude changes along such
109 paths, thus increasing the sensitivity, but also there is much better global averaging along such paths
110 thus giving more potential to measure long term effects, such as those due to global warming, with a
111 higher sensitivity.

112

113 **2. VLF Measurement Technique and Paths**

114 **2.1 The Portable VLF Loop Antenna and Receiver**

115 The phases and amplitudes of the VLF signals were measured both near and far from the
116 transmitters with a portable loop antenna with battery-powered circuitry. The phases were measured
117 (modulo half a cycle) relative to the 1-s pulses from a GPS receiver built in to the portable VLF
118 circuitry. The VLF signals came from NWC (North West Cape, Australia, 19.8 kHz), NPM (Oahu,
119 Hawaii, 21.4 kHz) or NLK (Seattle, 24.8 kHz) which, as for other US Navy VLF transmitters, are
120 modulated with 200 baud MSK. Details of the portable loop and its phase and amplitude measuring
121 techniques are given in *Thomson* [2010]. As previously, for measurements **at less than about 200**
122 **km** from the transmitters, the loop had extra resistance (typically $2 \times 750 \Omega$ or $2 \times 2k\Omega$) added in
123 series with it to reduce the gain. For all other measurements (far from the transmitters) this series
124 resistance was a nominal $2 \times 39 \Omega$. All phases (and amplitudes) reported here were either measured
125 with $2 \times 39 \Omega$ or adjusted to $2 \times 39 \Omega$ as in *Thomson* [2010]. The portable loop phase and amplitude
126 measurements used here were made on reasonably flat ground, away from significant hills, with most
127 being made in public parks or by the sides of (minor) roads. Care, as always, was needed to keep
128 sufficiently away from (buried/overhead) power lines and the like, particularly checking that
129 measurements were self-consistent over distances of at least a few tens of meters and from one

130 (nearby) site to the next. Some sites tried needed to be rejected but most, provided certain parts were
131 avoided, proved satisfactory and convenient.

132

133 **2.2 The Fixed VLF Recorders**

134 NWC, NPM and NLK, like other US Navy VLF transmitters, typically have very good phase and
135 amplitude stability. However, as with the other US transmitters, they normally go off-air once a week
136 for 6-8 hours for maintenance. On return to air, the phase is still normally stable but the value of the
137 phase (relative to GPS or UTC) is often not preserved. In addition, in the course of a typical week,
138 there may be some gradual phase drift or a small number of additional times when there are random
139 phase jumps. For meaningful phase comparisons, it was thus very desirable to have a fixed recorder
140 continuously recording while the portable measurements were being made. This was not convenient
141 to do locally in Australia, Hawaii or Seattle but was done near Dunedin, N.Z., where the signal-to-
142 noise ratio is still very good for NWC, NPM and NLK. The two recorders used, for both phase and
143 amplitude, were softPALs [Dowden and Adams, 2008] using two independent VLF receivers and
144 antennas (one loop and one vertical electric field) and GPS 1-s pulses as their phase references. These
145 recorders are part of the Antarctic-Arctic Radiation-belt Dynamic Deposition VLF Atmospheric
146 Research Consortium (AARDDVARK) [Clilverd et al., 2009]
147 (http://www.physics.otago.ac.nz/space/AARDDVARK_homepage.htm). Because of the stability of
148 the (daytime) propagation this provided a satisfactory method of recording, and compensating for,
149 transmitter phase drifts (or jumps).

150

151 **2.3 The Paths**

152 Figure 1 shows the locations of the NWC, NPM and NLK transmitters (diamonds), the principal
153 receiving locations (circles) and the great circle propagation paths (GCPs) which, as can be seen, are
154 mainly over the sea. The direction of propagation for each path is indicated by an arrow on its GCP.

155

156 **3. NWC to Tumwater (near Seattle)**

157 **3.1 Measurements of NWC at Tumwater**

158 Around 20 sets of portable loop phase and amplitude measurements of NWC signals were made in
159 and around Tumwater, WA (near Seattle, USA), over the 5 days 5-9 August 2008. Nearly all the
160 measurements were made **during the period ~0000-0230 UT, i.e., within ~2 hours of midday for**
161 **the path midpoint of the NWC-Tumwater path.** Five sites were used, mainly in public parks,
162 within ~2-12 km of each other. All the phase measurements were entered into an (Excel) spreadsheet
163 together with the site locations measured by a portable GPS receiver and later checked against
164 Google Earth. The spreadsheet was used to adjust the measured phase delays for the different ranges
165 from the transmitter (1.0 μs per 300 m) to allow comparison of sites. All the chosen sites gave
166 satisfactory results: on each of the 5 days the deviation from the mean phase of the (typically) 4 sites
167 used that day was $\sim\pm 0.5 \mu\text{s}$ (maximum $\sim\pm 0.8 \mu\text{s}$). The results from the site in Pioneer Park,
168 Tumwater, looked to be the most representative and reliable, and are shown in Table 1.

169 As previously [Thomson, 2010], all phase and amplitude measurements were taken in pairs: first
170 with the loop pointing directly ‘towards’ the transmitter and then, after rotation by 180° about the
171 vertical, pointing directly ‘away’ from the transmitter, thus reversing the phase of the magnetic field
172 but not the phase of any (unintentional residual) electric field. The two resulting amplitude
173 measurements in each pair seldom differed by more than $\sim 0.3 \text{ dB}$, usually less; similarly, the two
174 resulting phase measurements in each pair seldom differed by more than $\sim 0.5 \mu\text{s}$, usually less. For
175 each day, the table shows the average of the two 180° loop orientations for each of the two (sideband)
176 frequencies.

177 The second last column of the table shows the phase of NWC recorded at Dunedin as shown in
178 Figure 2a. The last column shows the Dunedin phase (in degrees) adjusted in line with the phases of
179 NWC observed at Tumwater as shown in columns 3 and 4. For example, the mean Tumwater phase

180 on 5 Aug 08 was $(21.3 + 19.8)/2 \mu\text{s} = 20.55 \mu\text{s}$, while on 7 Aug 08 it was $(17.0 + 15.6)/2 \mu\text{s} = 16.3$
 181 μs . This (apparent) decrease in phase delay of $20.55 - 16.3 \mu\text{s} = 4.25 \mu\text{s}$ from 5 to 7 Aug 08 is
 182 equivalent to an increase of the phase angle by $4.25 \times 10^{-6} \times 19800 \times 360^\circ = 30^\circ$; thus the “deg adj.”
 183 for 7 Aug 2008 relative to 5 Aug is $126^\circ - 30^\circ = 96^\circ$ as shown. From this last column of Table 1, it
 184 can be seen that the range of scatter for the measured phases for the (14.2 Mm) NWC to Tumwater
 185 path (relative to the NWC-Dunedin phases) is 18° or $\sim\pm 9^\circ$ from the mean, implying a likely random
 186 error of $\sim\pm 4^\circ$ for the mean of the NWC phase at Tumwater measured over the 5 days, 5-9 August
 187 2008.

188

189 **3.2 Observations and Modeling: NWC to Tumwater**

190 In a very similar manner to Table 1 here, Table 1 of *Thomson* [2010] showed the phases of NWC
 191 measured with the same portable loop system at Onslow, Western Australia, ~ 100 km ENE over the
 192 sea from NWC for the 3 days 21-23 October 2009. From these two tables, the mean Onslow and
 193 Tumwater phases ($19.3 \mu\text{s}$ and $17.6 \mu\text{s}$) and their corresponding Dunedin phases (-26° and 117°) were
 194 then used, in Table 2 here, to find the observed phase delay difference between Onslow and
 195 Tumwater. This, of course, required correcting for the phase changes at NWC (as measured at
 196 Dunedin) between the times of the Onslow and Tumwater measurements as shown in Table 2.

197 This delay difference (between Onslow and Tumwater) can be thought of as consisting of two parts:
 198 the free space part along the surface of the Earth and the ionospherically reflected part. Indeed
 199 programs such as ModeFinder and LWPC output their phases relative to the free-space delay. Table 3
 200 shows the locations of NWC and the principal sites used in each of Tumwater and Onslow (using
 201 Google Earth and a portable GPS receiver). The distances in rows 2 & 3 were calculated using the
 202 Vincenty algorithm [*Vincenty*, 1975; www.ngs.noaa.gov/cgi-bin/Inv_Fwd/inverse2.prl;
 203 www.ga.gov.au/geodesy/datums/vincenty_inverse.jsp] and from these the delays were found using
 204 the (exact) speed of light, $c = 299.792458 \text{ m}/\mu\text{s}$. The difference between the NWC-Tumwater and

205 NWC-Onslow delays, 47148.30 μs , was then reduced by an integral number of half cycles: 47148.30
 206 $- 1867 \times 0.5/0.0198 \mu\text{s} = 1.84 \mu\text{s}$, to allow for the phase measuring half-cycle ambiguity. This free
 207 space delay, modulo half-a-cycle, was then subtracted from the observed delay giving the waveguide
 208 part of the delay difference between Onslow and Tumwater, $18.4 - 1.84 \mu\text{s} = 16.6 \mu\text{s} \equiv 118^\circ$, which
 209 was then subtracted from the 128° calculated by LWPC (using $H' = 71.7 \text{ km}$ and $\beta = 0.43 \text{ km}^{-1}$) for
 210 the phase of NWC at Onslow in early August giving 10° , or equivalently $10^\circ - 180^\circ = -170^\circ$ (due to the
 211 half cycle ambiguity) as a preliminary value for the ‘observed’ phase at Tumwater shown in Figure
 212 2c. This preliminary phase value needs some seasonal refinement because of the different time of
 213 year that the measurements were made; the phases of NWC measured at Onslow (near NWC) during
 214 late October 2009 need to be adjusted to early August 2008 (when the Tumwater phases were
 215 measured) using NWC phases measured in Dunedin because this (5.7 Mm) NWC-Dunedin path will
 216 have undergone some seasonal changes in its phase delay in the 2.5 months between early August and
 217 late October. (The solar cycle changes will be minimal because both 2008 and 2009 were at solar
 218 minimum.)

219 Fortunately these seasonal phase changes for the NWC-Dunedin path over these 2.5 months can be
 220 fairly readily estimated. There are two principal effects. The first is changing H' and β , due to
 221 changing solar zenith angle over the period, the values for which were taken from *McRae and*
 222 *Thomson* [2000] and used in LWPC showing that a phase advance of 20° at Dunedin would be
 223 expected from early August to late October (mainly due to the decreasing solar zenith angle allowing
 224 the **Sun's** Lyman- α to penetrate deeper and so lower H'). The second effect is due to the warming of
 225 the neutral atmosphere as the southern-hemisphere season advances from winter towards summer,
 226 resulting in the height of a fixed atmospheric density (say 10^{21} m^{-3}) increasing and so H' increasing by
 227 the same amount. Neutral number density height profiles (for $[\text{N}_2]$) were found from the MSIS-E-90
 228 atmosphere model (http://omniweb.gsfc.nasa.gov/vitmo/msis_vitmo.html), around 70 km altitude in
 229 early August and late October from which it was found that H' increased, due to this warming effect,

230 by an average of ~ 1.35 km over the length of the NWC-Dunedin path during this period (see Figure
 231 3, discussed later). Using LWPC to model the effect of this 1.35 km height increase (without change
 232 in β) shows the phase at Dunedin would decrease by 22° due to this effect alone. The combination of
 233 these two effects means that phases in Dunedin in late October are to be expected to be just $22^\circ - 20^\circ$
 234 $= 2^\circ$ lower than in early August (for constant phase at NWC). A similar calculation shows the phase
 235 at Onslow would be $\sim 3^\circ$ higher in late October than in early August, due to these same two effects.
 236 This results in the preliminary -170° for the 'observed' phase at Tumwater found above becoming
 237 $-170^\circ + 2^\circ + 3^\circ = -165^\circ$.

238 The phase of NWC at Onslow was also measured on the 3 days 26-28 June 2008 [Thomson, 2010]
 239 just ~ 6 weeks before the Tumwater measurements, while recordings were being made in Dunedin.
 240 These June measurements have the advantage over the October Onslow measurements used above in
 241 that the predictable changes in the propagation (phase) on the NWC-Dunedin path over this winter 6
 242 weeks due to solar zenith angle (LWPC: $\sim 4^\circ$) and neutral temperature (MSIS-E-90: $\sim 5^\circ$) are much
 243 less than for the 2.5 months between August and October (20° and 22° respectively from above).
 244 Unfortunately the NWC-Dunedin propagation path was less stable 26-28 June 2008 (than in, say,
 245 October) as is not unusual in mid-winter. The phase angles at Dunedin over the 3 measurement days
 246 in June (when the phase of NWC itself was very stable) covered a range of $28^\circ (\pm 14^\circ)$ as compared
 247 with a range of only 2° in October (relative to a fixed phase at NWC or Onslow). Using the same
 248 process for adjusting the June Onslow phases from Dunedin recordings (not shown here) and with the
 249 same method of propagation corrections as for October (but now for June), the 'observed' phase at
 250 Tumwater in early August 2008 was estimated to be $-164^\circ + 4^\circ - 5^\circ = -165^\circ$, essentially the same as
 251 was obtained above by adjusting from the October measurements. Hence this -165° is shown in
 252 Figure 2c as the (final) 'observed' phase of NWC at Tumwater for comparison with modeling. The
 253 error in the mean 'observed' phase via the June Onslow phases will be largely due to the NWC-
 254 Dunedin propagation uncertainties and so $\sim \pm 10^\circ$ (i.e. somewhat less than the $\pm 14^\circ$ total measurement

255 range noted above) while the error in the mean via the October Onslow phases will be largely due to
 256 uncertainties in the NWC-Dunedin propagation changes between early August and late October,
 257 probably $\sim\pm 7^\circ$. Hence the error in the (final) 'observed' phase of NWC at Tumwater of -165° can be
 258 estimated to be $\sim\pm 6^\circ$.

259 The mean amplitude of the NWC signal measured at the Tumwater sites (14.2 Mm from NWC) at
 260 **mid-path midday (i.e. midday at the path midpoint)** on the five measurements days, 5-9 August
 261 2008, was $458 \mu\text{V/m} \equiv 53.2 \text{ dB}$ above $1 \mu\text{V/m}$. Virtually all of the measurements were within $\pm 1 \text{ dB}$
 262 of this value. (As can be seen in Figure 2b, NWC's amplitude at Dunedin was steady during this
 263 time.) There was significant atmospheric noise near Tumwater but the overall error in the mean
 264 amplitude at Tumwater is likely to be **less than approximately** $\pm 0.7 \text{ dB}$. The mean amplitude of the
 265 NWC signal measured at Onslow, 21-23 October 2009, was $99.7 \text{ dB} > 1 \mu\text{V/m}$ which indicates that
 266 NWC was radiating about 0.3 dB below 1 MW [Thomson, 2010]. (The same radiated power was also
 267 obtained from portable loop measurements in Onslow, 26-28 June 2008.) In Figure 2d the LWPC
 268 calculated amplitudes for the various values of H' and β are for a radiated power of 1 MW (being a
 269 convenient normalized value) but, to compensate for the apparently 0.3 dB lower radiated power, the
 270 'observed' amplitude is shown as $53.2+0.3 = 53.5 \text{ dB} > 1 \mu\text{V/m}$ (being the amplitude which would
 271 have been observed at Tumwater had NWC been radiating a full 1 MW).

272 It can thus be seen from the comparison between calculations and observations for the 14.2 Mm
 273 path NWC to Tumwater, in Figures 2c and 2d, that the best fit is for an ionosphere with $H' = 71.1 \text{ km}$
 274 and $\beta = 0.42 \text{ km}^{-1}$ averaged along this solar minimum path.

275

276 **3.3 Comparison with earlier Measurements and Modeling**

277 These average observed values of $H' = 71.1 \text{ km}$ and $\beta = 0.42 \text{ km}^{-1}$ for the long NWC-Tumwater
 278 path can usefully be compared with the values $H' = 70.5 \text{ km}$ and $\beta = 0.47 \text{ km}^{-1}$ for the short (300 km)
 279 low-latitude ($\sim 30^\circ$ geomagnetic) NWC-Karratha path (for near overhead **Sun**) [Thomson, 2010] and

280 the values $H' = 71.8$ km and $\beta = 0.34$ km⁻¹ for the short (360 km) high-midlatitude (~53.5°
 281 geomagnetic) NAA-PEI path (for near overhead **Sun**) [Thomson *et al.*, 2011]. The latter paper also
 282 gives a graph of β versus geomagnetic latitude interpolated using the known latitudinal variation of
 283 galactic cosmic ray fluxes. From this graph it can be seen that $\beta \approx 0.485$ km⁻¹ for the first 2/3 of the
 284 NWC-Tumwater path (~±30° geomagnetic) while the latter 1/3 (at the Tumwater/Seattle end) would
 285 have β varying between 0.47 and 0.34 km⁻¹ probably averaging about 0.41 km⁻¹ thus implying an
 286 average β for the path of $0.485 \times 2/3 + 0.41 \times 1/3 = 0.46$ km⁻¹ for midday **Sun** at all points along the
 287 path. By using the plot of β versus solar zenith angle given (from observations) by *McRae and*
 288 *Thomson* [2000] it can readily be estimated that the average value of β along the path will be lower
 289 by about 0.04 km⁻¹ (due to the higher solar zenith angles near the NWC and Tumwater ends of the
 290 path, even at **mid-path midday**) and so, based on the recent short-path results above, the expected
 291 average β would be $0.46 - 0.04 = 0.42$ km⁻¹ in close agreement with the direct results presented in
 292 Figure 2d.

293 As noted earlier and in *Thomson et al.* [2011], the principal source of variation in height, H' , with
 294 latitude and season seems to result from the changes in height of a fixed number density (e.g. 10^{21}
 295 m⁻³) in the neutral atmosphere which, as mentioned earlier, can be obtained from the MSIS-E-90
 296 model. This model was used to find the neutral density (actually $[N_2]$, the number density of N_2) at a
 297 height equal to $H' = 70.5$ km at the latitude of the (300-km) NWC-Karratha path in late October
 298 (2009) when this value of $H' = 70.5$ km was measured. This number density was $[N_2] = 1.31 \times 10^{21}$
 299 m⁻³. MSIS-E-90 was then used to find the height of this value of $[N_2]$ at other latitudes and times thus
 300 giving reasonable estimates for the values of H' (for near overhead **Sun**) for those times and places.
 301 These values of H' , as a function of latitude for early August (2008), are shown in Figure 3, as black
 302 squares, together with the baseline result, $H = 70.5$ km, for late October (2009) near NWC (red
 303 diamond and horizontal red dotted line). Figure 3 implies (black dashed line) that for early August the
 304 (average) value of H' , between 22° S and 30° N (i.e the first ~60% of the NWC-Tumwater path) is

305 70.15 km. For the remaining 40% of the path (latitudes 30°-47° N) the average value of H' can be
306 seen to be $(70.3+71.1)/2 = 70.7$ km. Thus the average for the whole path from this plot in Figure 3
307 (i.e. for midday at all points on the path) is $H' = 0.6 \times 70.15 + 0.4 \times 70.7 = 70.37$ km. The actual
308 average H' for the path will be a little higher than this because, at **mid-path midday**, the NWC end
309 will have morning solar zenith angles while the Tumwater end will have afternoon solar zenith
310 angles. These increases of H' with solar zenith angle towards the ends of the (midday) path can be
311 found from the appropriate plot in *McRae and Thomson* [2000]. The average increase in H' for the
312 first $\sim 1/3$ of the path (the NWC or morning end) was thus found to be **~ 1.7 km**. From this same plot,
313 the last $\sim 1/3$ of the path (the Tumwater or afternoon end) would also look to have an increase of **~ 1.7**
314 **km**. However, this end is at a significantly higher geomagnetic latitude, and so has a much higher
315 proportion of its electron density from (zenith angle independent) cosmic rays than the *McRae and*
316 *Thomson* [2000] plot. Hence less variation in H' with solar zenith angle is to be expected at the
317 Tumwater end. A similar latitude situation exists for the path NAA to Cambridge for which
318 amplitude and phase plots were available [*Thomson et al.*, 2007] allowing LWPC to be used to find
319 changes in H' with solar zenith angle for this high-midlatitude path. It was thus found that the
320 average increase in H' for the afternoon $1/3$ of the NWC-Tumwater would likely be only ~ 0.8 km.
321 Hence the (final) value of H' , averaged along the NWC-Tumwater path (at *mid-path midday*), from
322 all these earlier observations would be $H' = 70.37 + (1.7 + 0.8)/3$ km = 71.2 km. This is only ~ 0.1 km
323 higher than the 71.1 km obtained from the present direct measurements on the NWC-Tumwater path
324 shown in Figure 2c which is thus very satisfactory. **These comparisons between the long-path**
325 **measurements of H' and β with the corresponding values from short-path results adjusted for**
326 **changing solar zenith angle and changing latitudes along the path are summarized in Table 4**
327 **(columns 2 and 3).**

328

329

330 **4. NPM (Hawaii) to Dunedin, NZ**

331 Measurements similar to those for the NWC to Tumwater path were also made for the ~8.1 Mm
 332 NPM to Dunedin path. The US Navy 21.4 kHz transmitter, NPM (on the Hawaiian Island of Oahu), is
 333 located at 21.4202° N, 158.1511° W. Phases and amplitudes of NPM were measured with the portable
 334 loop system at several suitable sites on the eastern side of the nearby island of Kauai, on the four days
 335 27, 28, 30 and 31 Oct 2009 (NPM was off-air for ~8 hours until ~02 UT on 29 Oct 2009). The prime
 336 receiving site there (which gave readings consistent with those at the other sites on Kauai) was in
 337 Lydgate Park located at 22.0385° N, 159.3362° W, which was thus 140.42 km from NPM (using the
 338 **Vincenty** algorithm). Phase and amplitude recordings of NPM were made at Dunedin (using softPAL
 339 recorders) before, during and after the Kauai measurements. Portable loop measurements of NPM's
 340 phase and amplitude were made at several sites in Dunedin (giving good mutual agreement) both
 341 before and after the Kauai measurements. The prime (reference) site in Dunedin was in Bayfield Park
 342 at 45.8938° S, 170.5236° E which, using the **Vincenty** algorithm, is thus 8098.08 km from NPM. The
 343 path difference between Bayfield Park and Lydgate Park was thus found to be $8098.08 - 140.42 \text{ km} =$
 344 7957.65 km which, using the (exact) speed of light, corresponds to a free-space delay of $26543.87 \mu\text{s}$
 345 which, modulo half-a-cycle of NPM's 21.4 kHz (i.e. $0.5/0.0214 \mu\text{s}$) becomes $1.81 \mu\text{s}$. The
 346 corresponding observed phase delay (from the portable loop phase measurements in Bayfield and
 347 Lydgate Parks) was found (in a similar manner to that for NWC and Tumwater in Section 3) to be 5.2
 348 μs , which means the "waveguide only" part of the delay was $5.2 - 1.8 \mu\text{s} = 3.4 \mu\text{s}$, **modulo a quarter**
 349 **of a period of 21.4 kHz, because the 21.4 kHz phase measurement is derived from the (portable**
 350 **loop) 21.35 kHz and 21.45 kHz sideband measurements, either or both of which have**
 351 **(independent) half-cycle ambiguities.** This "waveguide only" delay of $3.4 \mu\text{s} \equiv 26^\circ$ (modulo 90°) is
 352 then subtracted from the 127° phase found by LWPC (using $H' = 71.8 \text{ km}$ and $\beta = 0.44 \text{ km}^{-1}$) for NPM
 353 at Lydgate in October to get the $127^\circ - 26^\circ - 90^\circ = 11^\circ$ shown as the "observed" phase for NPM at

354 Dunedin in Figure 4 which also shows the LWPC-calculated phases for NPM at Dunedin for
 355 appropriate values of H' and β .

356 The mean amplitude of the NPM signal measured at the Dunedin sites (~ 8.1 Mm from NPM) at
 357 **mid-path midday** (~ 23 UT) in October/November 2009 was $460 \mu\text{V/m} \equiv 53.3$ dB above $1 \mu\text{V/m}$
 358 which is shown as the "observed" amplitude in Figure 4 for comparison with LWPC modeling.
 359 Virtually all of the measurements were within ± 0.7 dB of this value so that the error in the mean is
 360 likely to be $\sim \pm 0.5$ dB. On Kauai, ~ 140 km from NPM, the measured effective midday mean amplitude
 361 of the NPM signal, 27-31 Oct 2009, was 40.6 ± 2 mV/m $\equiv 92.2$ dB $> 1 \mu\text{V/m}$. Using LWPC, with an
 362 appropriate (midday, late October, 22° N) ionosphere, $H' = 71.8$ km, $\beta = 0.44 \text{ km}^{-1}$, on this NPM-
 363 Kauai path, gave the radiated power as 375 kW. This power was then used again in LWPC to calculate
 364 the expected amplitudes of NPM at Dunedin (8.1 Mm away) for appropriate values of H' and β giving
 365 the results shown in Figure 4.

366 From Figure 4 it can be seen that $H' = 70.8$ km and $\beta = 0.46 \text{ km}^{-1}$ give good fits to the observed
 367 phases and amplitudes for NPM-Dunedin. These average observed parameters for this fairly long path
 368 can again usefully be compared with the recent short-path parameters as was done for NWC-
 369 Tumwater in Section 3.3. **A summary is given in Table 4 (columns 4 and 5).** Because the NPM-
 370 Dunedin path is much shorter (8.1 Mm compared with 14.2 Mm) and more North-South (covering less
 371 local time), the variations in H' and β along the path are much smaller, but can be dealt with in a very
 372 similar manner. For the $\sim 70\%$ of the NPM-Dunedin path with low geomagnetic latitudes between 21°
 373 N and 30° S, the average β (as before) will be $\sim 0.485 \text{ km}^{-1}$, while for the remaining 30% of the path
 374 (the Dunedin end) the average β will be $(0.47+0.34)/2 \text{ km}^{-1} = 0.41 \text{ km}^{-1}$ giving $\beta = 0.7 \times 0.485 + 0.3 \times$
 375 $0.41 \text{ km}^{-1} = 0.463 \text{ km}^{-1}$ for the path average, for the **Sun** near the zenith all along the path. The effects
 376 of the actual higher solar zenith angles near the ends of the path (at **mid-path midday**) can be
 377 estimated, as before, from *McRae and Thomson* [2000], as reducing β by $\sim 0.006 \text{ km}^{-1}$ thus giving $\beta =$
 378 $0.463 - 0.006 \text{ km}^{-1} = \sim 0.46 \text{ km}^{-1}$ which agrees very well with the 0.46 km^{-1} directly measured here on

379 the long NPM-Dunedin path. From the green lines in Figure 3, the average value of H' (for near
 380 overhead **Sun**) for the NPM-Dunedin path (21° N to 46° S, late October) is 70.55 km. From *McRae*
 381 *and Thomson* [2000], the small increases in H' near the ends of the path (due to the higher solar zenith
 382 angles there at **mid-path midday**) can be estimated (as for the NWC-Tumwater path in section 3).
 383 This resulted in $H' = 70.55 + 0.25 \text{ km} = 70.8 \text{ km}$ which again agrees very well with the 70.8 km found
 384 here from direct observations on the long NPM-Dunedin path.

385

386 **5. NWC to Kauai, Hawaii**

387 Measurements similar to those for the NWC-Tumwater path were also made for the ~10.6 Mm
 388 path NWC to Kauai. Phases and amplitudes of NWC were measured with the portable loop system at
 389 several suitable sites on the eastern side of the Island of Kauai on five days, 27-31 Oct 2009. The
 390 prime receiving site there (which gave readings consistent with those at the other sites on Kauai) was
 391 the same site in Lydgate Park as used for NPM (section 4); this site was 10560.92 km from NWC
 392 (again making use of the **Vincenty** algorithm). From section 3.2, Table 2, the distance from NWC to
 393 (the prime site in) Onslow was 100.16 km so that the Lydgate-Onslow path difference is $10560.92 -$
 394 $100.16 \text{ km} = 10460.76 \text{ km}$ which, using the (exact) speed of light gives the free-space delay difference
 395 as $34893.35 \mu\text{s}$ which, in turn, **modulo half-a-period of NWC's 19.8 kHz (0.5/0.0198 μs), becomes**
 396 **19.6 μs** . As mentioned previously (section 3.2 and *Thomson* [2010]), the phase of NWC was measured
 397 with the portable loop system at Onslow, 21-23 Oct 2009. Phase (and amplitude) recordings of NWC
 398 were also made at Dunedin (using softPAL recorders) before, during and after these Onslow and Kauai
 399 measurements to monitor and correct for any phase changes at NWC during this period. The NWC-
 400 Dunedin propagation path was, as usual, very stable during this late spring, solar minimum period
 401 making the Dunedin recorders very effective in monitoring the phase changes of the NWC transmitter.
 402 With the help of these Dunedin recordings using a very similar procedure to that for the NWC-
 403 Tumwater path (in section 3.2), the portable loop measurements at Lydgate Park and Onslow, gave the

404 observed Lydgate-Onslow phase difference (modulo half-a-cycle of NWC) as 19.6 μs . **Subtracting**
 405 **the calculated free-space delay of 19.6 μs (from above) from this observed 19.6 μs then**
 406 **gave 0 $\mu\text{s} \equiv 0^\circ$ or, modulo half-a-cycle, 180° , for the waveguide-only part of the Onslow-Lydgate**
 407 **delay.** Subtracting this 180° from the 131° calculated by LWPC (using $H' = 70.5 \text{ km}$, $\beta = 0.47 \text{ km}^{-1}$)
 408 for the phase of NWC at Onslow in late October gave -49° which is thus shown as the "observed"
 409 phase of NWC at Lydgate Park in Figure 5 where it is compared with the LWPC-calculated NWC
 410 phases at Lydgate Park using suitable values of H' and β .

411 The mean amplitude of the NWC signal measured at the Kauai sites ($\sim 10.6 \text{ Mm}$ from NWC) at
 412 **mid-path midday** ($\sim 01 \text{ UT}$) 27-31 October 2009 was $590 \mu\text{V/m} \equiv 55.4 \text{ dB}$ above $1 \mu\text{V/m}$. Figure 5
 413 shows the LWPC-calculated amplitudes at Kauai for NWC radiating 1 MW. As noted in section 3.2,
 414 the Onslow portable loop measurements were more consistent with NWC radiating $\sim 0.3 \text{ dB}$ below 1
 415 MW. The "observed" amplitude for NWC is thus shown in Figure 5 as $55.4 + 0.3 \text{ dB} = 55.7 \text{ dB}$ (as
 416 would have been observed if NWC had been radiating a full 1 MW).

417 From Figure 5 it can be seen that $H' = 71.0 \text{ km}$ and $\beta = 0.46 \text{ km}^{-1}$ give good fits to the observed
 418 phases and amplitudes for NWC-Kauai. These average observed parameters for this long path can
 419 again usefully be compared with the recent short-path values as was done for NWC-Tumwater in
 420 Section 3.3 and NPM-Dunedin in section 4. **Again a summary is given in Table 4 (columns 6 and**
 421 **7).** Because all of the NWC-Kauai path is at low geomagnetic latitudes (between 21° N and 30° S), the
 422 average β (as before) would be $\sim 0.485 \text{ km}^{-1}$ if the **Sun** were near overhead at all points along the (10.6
 423 Mm) path. The effects of the higher solar zenith angles near the ends of the path (at **mid-path**
 424 **midday**) can be estimated, as previously, from *McRae and Thomson* [2000], as reducing β by ~ 0.016
 425 km^{-1} thus giving $\beta = 0.485 - 0.016 \text{ km}^{-1} = \sim 0.47 \text{ km}^{-1}$ which agrees quite well with the 0.46 km^{-1} found
 426 above from the directly-measured, long NWC-Kauai path. From Figure 3, the average value of H' (for
 427 near overhead **Sun**) for the NWC-Kauai path (22° S to 22° N) is 70.6 km in late October. From
 428 *McRae and Thomson* [2000], the small increases in H' near the ends of the path (due to the higher solar

429 zenith angles there at **mid-path midday**) can be estimated (as for the NWC-Tumwater path in section
430 3). This resulted in $H' = 70.6 + 0.55 \text{ km} = 71.15 \text{ km}$ which again agrees quite well with the 71.0 km
431 found here from direct observations on the long NWC-Kauai path.

432 However, unlike the two previously discussed paths here (NWC-Tumwater and NPM-Dunedin),
433 as can be seen in Figure 1, this NWC-Kauai path passes over significant amounts of land; ~1.8 Mm of
434 this ~10.6 Mm NWC-Kauai path is over northern Australia where the (LWPC-built-in) ground
435 conductivity is quite low, $\sim 1 \times 10^{-3} \text{ S/m}$. As an example of the sensitivity due to ground conductivity,
436 when the LWPC calculation for NWC-Kauai (with $H' = 71.0 \text{ km}$ and $\beta = 0.46 \text{ km}^{-1}$) was repeated
437 with an all-sea conductivity, the calculated amplitude at Kauai increased by ~4.3 dB and the phase
438 advanced by $\sim 27^\circ$. Given the considerable uncertainties in the ground conductivities, the agreement
439 between the estimated and observed values of β and H' for this NWC-Kauai path is remarkably good.

440

441 **6. NLK (Seattle) to Dunedin, NZ**

442 Similar measurements for the long ~12.3 Mm, nearly all-sea path NLK (24.8 kHz, Seattle) to
443 Dunedin were also made. NLK is located at 48.2036° N , 121.9171° W (from Google Earth). Phases
444 and amplitudes of NLK were measured with the portable loop system at several suitable sites ~150 km
445 south of Seattle in the vicinity of Olympia/Tumwater WA, USA, 5-9 August 2008. The prime
446 receiving site there (which gave readings consistent with other nearby sites within a few km) was the
447 same site in Pioneer Park (Tumwater) as used in Section 3.2 (Table 3), which is 152.60 km from NLK
448 (using the **Vincenty** algorithm). Phase and amplitude recordings of NLK were made at Dunedin (using
449 softPAL recorders) before, during and after the Tumwater measurements. Portable loop measurements
450 of the phase and amplitude of NLK were made at several sites in Dunedin (giving quite good mutual
451 agreement) both before and after the Tumwater measurements. The prime (reference) site in Dunedin
452 was again in Bayfield Park (as used in Section 4) which, using the **Vincenty** algorithm, is 12315.74
453 km from NLK. The path difference between Bayfield Park and Pioneer Park was thus found to be

454 12315.74 - 152.60 **km** = 12163.15 km which, using the (exact) speed of light, corresponds to a free-
 455 space delay of 40571.89 μs which, modulo half-a-cycle of NLK's 24.8 kHz (i.e. $0.5/0.0248 \mu\text{s}$)
 456 becomes 7.37 μs . The corresponding observed phase delay (from the portable loop measurements in
 457 Bayfield and Pioneer Parks) was found (in a similar manner to that for NWC and Tumwater in Section
 458 3.2) to be 17.7 μs , which means the "waveguide only" part of the delay was $17.7 - 7.4 \mu\text{s} = 10.3 \mu\text{s} \equiv$
 459 92° . This 92° was then subtracted from the 127° calculated by LWPC (using $H' = 71.7 \text{ km}$, $\beta = 0.33$
 460 km^{-1}) for the phase of NLK at Tumwater in early August, giving 35° , or equivalently $35^\circ - 180^\circ =$
 461 -145° (due to the half cycle ambiguity), for the 'observed' phase at Tumwater shown in Figure 6a
 462 which also shows the LWPC-calculated phases for NLK at Dunedin for appropriate values of H' and
 463 β .

464 The mean amplitude of the NLK signal measured at the Dunedin sites ($\sim 12.3 \text{ Mm}$ from NLK) at
 465 **mid-path midday** ($\sim 23 \text{ UT}$) in July/August 2008 was $65 \mu\text{V/m} \equiv 36.3 \text{ dB}$ above $1 \mu\text{V/m}$; this is
 466 shown as the "observed" amplitude in Figure 6b for comparison with the LWPC modeling. Nearly all
 467 of these measurements were within $\pm 1 \text{ dB}$ of this value so that the error in the mean is likely to be
 468 $\sim \pm 0.7 \text{ dB}$. At Tumwater, $\sim 153 \text{ km}$ from NLK, the measured effective midday mean amplitude of the
 469 NLK signal, 5-9 Aug 2008, was $31.7 \pm 2 \text{ mV/m} \equiv 90.0 \text{ dB} > 1 \mu\text{V/m}$. Using LWPC, with an
 470 appropriate (midday, early August, 47.5° N) ionosphere, $H' = 71.7 \text{ km}$, $\beta = 0.33 \text{ km}^{-1}$, on this NLK-to-
 471 Tumwater path, gave the radiated power as 290 kW . This power was then used again in LWPC to
 472 calculate the expected amplitudes of NLK at Dunedin (12.3 Mm away) for appropriate values of H'
 473 and β giving the results shown in Figure 6b for comparison with the observed amplitude.

474 As for the other three long paths discussed here (sections 3, 4 and 5), H' and β for this long path
 475 can again usefully be estimated from the recently measured short-path parameters. **Again a summary**
 476 **is given in Table 4 (columns 8 and 9)**. From Figure 3, the average value of H' (for near overhead
 477 **Sun**) for the NLK-Dunedin path (48° N to 46° S , early August) is 70.0 km . From *McRae and*
 478 *Thomson* [2000], the increases in H' near the ends of the path (due to the higher solar zenith angles

479 there at **mid-path midday**) can be estimated (as for the NWC-Tumwater path in section 3.3). This
 480 resulted in $H' = 70.0 + 0.8 \text{ km} = 70.8 \text{ km}$. For the ~60% of this NLK-Dunedin path with low
 481 geomagnetic latitudes between 30° N and 30° S , the average β (as before) will be $\sim 0.485 \text{ km}^{-1}$, while
 482 for the remaining 40% of the path (the Seattle and Dunedin ends) the average β will be $\sim (0.47+0.34)/2$
 483 $\text{km}^{-1} = 0.41 \text{ km}^{-1}$ giving $\beta = 0.6 \times 0.485 + 0.4 \times 0.41 \text{ km}^{-1} = 0.455 \text{ km}^{-1}$, for the **Sun** near the zenith all
 484 along the path. The effects on β of the actual higher solar zenith angles near the ends of the path (at
 485 **mid-path midday**) can be estimated, as before, from *McRae and Thomson* [2000], as reducing β by
 486 $\sim 0.02 \text{ km}^{-1}$ giving $\beta = 0.455 - 0.02 \text{ km}^{-1} = \sim 0.435 \text{ km}^{-1}$ which does not agree very well with the ~ 0.38
 487 km^{-1} indicated by the direct observations in Figure 6b. Indeed, as can be seen in Figure 6b, it appears
 488 that $\beta = 0.435 \text{ km}^{-1}$ would give an observed amplitude at Dunedin of $38.7 \text{ dB} > 1 \mu\text{V/m}$ whereas the
 489 actual portable loop observations gave 36.3 dB at Dunedin.

490 This apparent discrepancy appears to be a result of the effective radiated power from NLK being
 491 somewhat direction dependent. NLK is unusual in that, instead of using very tall towers (400-500 m
 492 high) on flat ground to make the antenna high enough to get a reasonable radiation efficiency, it has
 493 wires strung between mountain ridges across a valley with the radiating current coming up to these in
 494 a cable from the transmitter on the valley floor below [e.g., *Watt*, 1967]. As well as the NLK
 495 amplitude measurements near Tumwater ($\sim 153 \text{ km}$ SSW of NLK), additional amplitude measurements
 496 were made over a much greater land area and range of directions \sim SW of NLK (the closest to NLK
 497 being at Dosewallips State Park, 93 km from NLK, while the furthest was at Westport on the west
 498 coast, 220 km from NLK). A range-corrected plot of these measured amplitudes of NLK as a function
 499 of azimuth (degrees east of north from NLK) is shown in Figure 6c, where it can be seen that the
 500 amplitudes measured at sites in the direction of Dunedin are $\sim 3 \text{ dB}$ lower than those measured at sites
 501 near Tumwater. As the amplitudes at Tumwater were used to determine the radiated power of 290 kW
 502 used for NLK in calculating the amplitudes at Dunedin in Figure 6b, it seems that the low amplitudes
 503 ($\sim 36 \text{ dB}$) measured at Dunedin may well be due to the lower radiated power in this direction. Thus

504 quite likely the value of $\beta = 0.435 \text{ km}^{-1}$ estimated from the earlier short path measurements will be
505 more appropriate than the (radiation-direction compromised) value, $\beta = 0.38 \text{ km}^{-1}$, from amplitude
506 comparisons in Figure 6b. Indeed, if $\beta = 0.435 \text{ km}^{-1}$ is used in the NLK-Dunedin phase plot in
507 Figure 6a then this gives $H' = 70.9 \text{ km}$ in close agreement with the $H' = 70.8 \text{ km}$ estimated above from
508 the short path parameters.

509

510 **7. Discussion, Summary and Conclusions**

511 Phases and amplitudes of suitable VLF signals were measured using a portable loop system
512 referenced to GPS 1-second pulses. Observations of the midday VLF radio phase changes and
513 amplitude attenuations along four long, mainly all-sea paths have been presented here: NWC (N.W.
514 Australia) to Seattle (14.2 Mm), NPM (Hawaii) to Dunedin, N.Z. (8.1 Mm), NWC to Kauai, Hawaii
515 (10.6 Mm) and NLK (Seattle) to Dunedin, N.Z. (12.3 Mm). Average values of the height, H' , and
516 sharpness, β , of the D-region of the ionosphere along each path were then determined by modeling
517 with the waveguide code, LWPC, so that the modeled phases and amplitudes agreed with those
518 observed.

519 These resulting average values of H' and β for each of the four long paths were then compared with
520 recent short ($\sim 300 \text{ km}$) path measurements of H' and β at (i) a low geomagnetic latitude and (ii) a
521 mid-high geomagnetic latitude. The interpolation of β with geomagnetic latitude along the paths was
522 obtained from these by using the known variation of cosmic ray flux with geomagnetic latitude. The
523 variations (interpolations) of H' with season and geographic latitude along the paths were obtained
524 from the short-path observations by extending them with the MSIS-E-90 neutral atmosphere model.
525 Small additional variations of H' and β due to changes in solar zenith angle near the ends of the paths
526 were estimated from the observations of *McRae and Thomson* [2000]. **Han and Cummer [2010]**
527 **measured H' near Duke University, $\sim 37^\circ \text{ N}$ geographic, in summer using natural lightning.**

528 **Their values of H' for near overhead Sun are typically in the range 71-72 km, averaging ~ 71.5**
 529 **km which is a little greater than the $H' = 70.9$ km from MSIS-E-90 and Figure 3 in June/July**
 530 **here but is quite likely just within the combined experimental errors of the two methods.**

531 For the NPM-Dunedin path (8.1 Mm), both the direct long-path method and the
 532 interpolated/extrapolated short-path method gave essentially the same results, $H' = 70.8$ km and $\beta =$
 533 0.46 km^{-1} . The very long (14.2 Mm) NWC-Seattle path gave the same (average) value of $\beta = 0.42 \text{ km}^{-1}$
 534 with both short- and long-path methods **while the $H' = 71.1$ km obtained from the long path was**
 535 **only marginally lower, by ~ 0.1 km (~ 100 m), than the $H' = 71.2$ km obtained using the short-**
 536 **path method.** A similar small height difference was also seen on the (10.6 Mm) NWC-Kauai path (H'
 537 $= 71.0$ and 71.15 km) but this is of even less significance because of the uncertainty in the (low)
 538 conductivity of the ~ 1.8 Mm of (Australian) ground on this path. Similarly the difference between the
 539 short-path and long-path values of β for this path (0.47 and 0.46 km^{-1}) is also probably not enough to
 540 be significant for the same uncertain ground conductivity reason. For the (12.3 Mm) NLK-Dunedin
 541 path, in contrast with the two NWC paths, the $H' = 70.9$ km from the long-path method was slightly
 542 higher than the $H' = 70.8$ km from the short-path method. The uncertainties for β on this path, due to
 543 NLK and some of its nearby measurements being undertaken in mountainous terrain, mean that this
 544 small difference in H' (0.1 km) is probably not significant for this path.

545 Overall, the agreement between the short-path and long-path observations is remarkably good with
 546 maximum differences of only ~ 0.15 km in height, H' , and 0.01 km^{-1} in sharpness, β . This is suggestive
 547 that the errors ($\sim \pm 0.5$ km for H' and $\sim \pm 0.03 \text{ km}^{-1}$ for β) for the short-path measurements reported in
 548 Thomson [2010] and Thomson *et al.* [2011] may have been a little conservative. Because both the
 549 transmitters and the measurement sites there were on land, and the paths were short (~ 300 km), there
 550 was a concern that the resulting proportion of low-conducting land and coastal boundaries (though
 551 appreciably less than for 50% of the paths) might have been having more effect than hoped. Any
 552 concern that this might have been a difficulty is now markedly reduced. The use of the MSIS model

553 for estimating changes in H' with season and latitude (but not with solar zenith angle) effectively
554 assumes there are no relevant neutral atmosphere composition changes (near 70 km altitude) with
555 season and latitude. For the cosmic rays, which ionize all atmospheric constituents, this is very likely
556 to be the case. However, for the minor but important constituent NO (ionized by Lyman- α) this is
557 less certain. None-the-less the apparent agreement resulting from using MSIS with this assumption
558 may well be implying the proportion of NO in the neutral atmosphere at heights near 70 km, at least
559 for the low and midlatitudes studied here, is fairly constant with season and latitude.

560 The validated, quiet-time, daytime modeling presented here provides an improved baseline for
561 measuring a wide variety of perturbations to the lower D-region and hence the Earth-ionosphere
562 waveguide. Such 'perturbations' include the transition from day to night (so improving nighttime
563 parameters [*McRae and Thomson, 2007, 2009*]), the effects of solar flares [e.g., *Thomson et al., 2005*]
564 and the effects of particle precipitation [e.g., *Rodger et al., 2007*], which can perturb the day or night
565 ionosphere.

566 Of course, although the modeling by LWPC seems very good, it will not be perfect; better modeling
567 will be found in the future. In particular, the representation of the electron density versus height
568 profile in the lower D-region by the two simple parameters, H' and β , though very good, is not likely
569 to be exact. However, the raw phase and amplitude measurements for the long paths measured here
570 are independent of the current modeling; these measurements could well be used in future, improved
571 modeling thus determining (retrospectively) improved values for the height (and sharpness) of the
572 lowest edge of the (D-region of) the Earth's ionosphere. The sensitivity of these long paths to these D-
573 region parameters is quite high; the error of $\pm 6^\circ$ in phase and ± 0.7 dB in amplitude estimated for the
574 14.2 Mm NWC-Tumwater path in section 3.2 corresponds (using Figure 2) to a sensitivity of less than
575 approximately ± 0.3 km in H' and $\sim \pm 0.01$ km $^{-1}$ in β . Also, as can be seen from Figure 2c, a (short
576 term) change in height, H' , (by say 1 km) while β remained constant, caused by (say) a simple vertical
577 height change in the neutral atmosphere alone (i.e., where the height of a fixed density, say 10^{21} m $^{-3}$,

578 changes by 1 km) will cause a phase change of $\sim 44^\circ$ per km. In such a (constant β) case, the long term
579 phase measurement sensitivity of $\pm 6^\circ$ corresponds to a height sensitivity of $\sim \pm 0.15$ km. For very short
580 term changes over a few hours, sensitivities as low as $\pm 4^\circ$ or ± 0.1 km = ± 100 m are likely for this very
581 long path (if β is constant). Over much longer times, when β may not be quite constant, and even
582 without improved modeling, future long-path measurements of H' and β can potentially measure
583 changes in height (of ~ 0.3 km) over time to a higher accuracy than current modeling can determine
584 absolute height. Clearly such potential observations could include determining H' changes over a
585 solar cycle. It might also be possible to use such height changes (averaged over individual solar
586 cycles), over even longer periods of time, to test theories of global warming and the corresponding
587 height-integrated atmospheric expansions and contractions below a specified height in the D-region,
588 such as H' , which corresponds effectively to the (fixed) density level down to which the external
589 ionizing radiations (Lyman- α , galactic cosmic rays etc) penetrate before they are absorbed or their
590 effects are overwhelmed by electron loss processes (attachment/recombination).

Acknowledgements

The authors are very grateful their colleague, David Hardisty, for his design, development and construction of the VLF phase meter. Thanks also to <http://rimmer/ngdc.noaa.gov> of NGDC, NOAA, US Dept. of Commerce, for the digital coastal outline used in Figure 1.

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Figure Captions

Figure 1. The transmitter sites (red diamonds), the receiver sites (blue circles) and the long paths across the Pacific Ocean used for the VLF phase and amplitude measurements.

Figure 2. (a) NWC phases and (b) amplitudes recorded at Dunedin, N.Z., while portable loop measurements of NWC were being made at Tumwater (near Seattle). Comparisons of observed midday (c) phases and (d) amplitudes (using NWC-Dunedin as reference) with modeling by LWPC for the NWC to Tumwater path.

Figure 3. Variations of H' (in km, for near overhead Sun) with latitude and season due to neutral atmosphere changes from the MSIS-E-90 atmospheric model. The plots are only slightly longitude dependent. The red diamond and its associated red dotted line are the reference height (H' as measured for the short, 300 km, NWC-Karratha path). The other dotted and dashed lines are used to aid in visualizing the averaging of H' along some of the long paths as discussed in the text.

Figure 4. Comparisons of observed midday phases and amplitudes with modeling for the NPM to Dunedin path.

Figure 5. Comparisons of observed midday phases and amplitudes with modeling for the NWC to Kauai, Hawaii, path.

Figure 6. Comparisons of observed midday (a) phases and (b) amplitudes with modeling for the NLK to Dunedin path plus (c) the observed directivity of NLK towards Dunedin and Tumwater.

Tables

Table 1. NWC Phases Measured at Tumwater, WA, and Dunedin^a

UT Date	UT	L (μs) ^b	H (μs) ^b	Dn (deg) ^b	adj. (deg)
5 Aug 08	0022	21.3	19.8	103	103
6 Aug 08	0240	21.3	19.8	90	90
7 Aug 08	0027	17.0	15.6	126	96
8 Aug 08	0018	8.7	6.7	197	106
9 Aug 08	0124	24.0	21.8	71	88

^aThe phase measurements at Tumwater were observed using $2 \times 39 \Omega$ and are in μs (L = 19.75 kHz, H = 19.85 kHz). Measurements at Dunedin are in degrees. The last column, 'adj.', illustrates the consistency (while NWC's phase drifts) by adjusting the 'Dn' phase in line with the Tumwater ' μs ' phase, as explained in the text.

^bMean for these ten Tumwater phases is 17.6 μs . Mean for these five phases at Dunedin is 117°.

Table 2. Observed NWC Phase Difference between Tumwater, WA, USA and Onslow, N.W. Australia^a

Observed	Phase (μs)	Dn (deg)
Tumwater $2 \times 39 \Omega$	17.6	117
Onslow $2 \times 39 \Omega$	19.3	-26
Onslow $2 \times 39 \Omega$	-0.8	117
Δ Phase (Tumwater-Onslow)	18.4	—

^aThe observed phase difference between Tumwater and Onslow (row 4), after correcting the measured Onslow phase (see text, shown here in row 2) for the NWC phase drift as measured at Dunedin (row 3) between the times of the Onslow and Tumwater (row 1) observations.

Table 3. Calculated Onslow-Tumwater Free-Space Delay Differences^a

Calculated Phases (μs)	Latitude (deg)	Longitude (deg E.)	Dist. (km)	Delay (μs)
NWC	-21.8163	114.1656		
Tumwater (Pioneer Park)	46.9970	-122.8843	14234.86	47482.4
Onslow	-21.6374	115.1146	100.16	334.1
Δ f: Tumwater – Onslow			14134.71	47148.3
Δ f: modulo half cycle				1.84
Δ o: observed (ex Table 2)				18.4
W/guide delay (Δ o – Δ f)				16.6

^aRows 1–4 show the locations with calculated distances and free space delays for NWC-Tumwater, NWC-Onslow and Onslow-Tumwater. Row 5 then shows the Onslow-Tumwater free-space delay difference modulo half a cycle of 19.8 kHz. This difference is then subtracted from the 18.4 μs observed delay (row 6), from Table 2, to give the waveguide only part of the delay as 16.6 μs (bottom row) which is equivalent to 118°. This observed 118° is then subtracted from the 128° calculated by LWPC for Onslow giving $10^\circ - 180^\circ = -170^\circ$ which is used, after small seasonal adjustments (see text), in Figure 2c as the 'observed' NWC phase at Tumwater.

Table 4. Comparison of measured long-path H' and β with values from previous measurements and available sources^a

Data Source (details given in text)	NWC- Tumwater (~14 Mm)	NWC- Tumwater (~14 Mm)	NPM- Dunedin (~8 Mm)	NPM- Dunedin (~8 Mm)	NWC- Kauai (~11 Mm)	NWC- Kauai (~11 Mm)	NLK- Dunedin (~12 Mm)	NLK- Dunedin (~12 Mm)
	H' (km)	β (km^{-1})	H' (km)	β (km^{-1})	H' (km)	β (km^{-1})	H' (km)	β (km^{-1})
Short paths - mid-day Sun all along the path	70.37	0.46	70.55	0.463	70.6	0.485	70.0	0.455
Solar zenith angle adjustment at ends of path	+0.83	-0.04	+0.25	-0.006	+0.55	-0.016	+0.8	-0.02
Results from combining the two rows above	71.2	0.42	70.8	0.46	71.15	0.47	70.8	0.435
Long-path measurements reported here	71.1	0.42	70.8	0.46	71.0	0.46	70.9	(0.38)

^aThe values of H' and β derived from the long-path VLF phase and amplitude measurements reported here are summarized in row 4 for each of the four paths. The corresponding values of H' and β derived from earlier measurements and sources, for constant (mid-day) solar zenith angle along the paths, are shown in row 1. Row 2 shows the adjustments needed for the row 1 values to allow for the higher solar zenith angles towards the ends of the paths. Row 3 combines rows 1 and 2 for comparisons with row 4. (See text for details.)











